

# Simulations / statistical mechanics / ...

## Topics

- very simple stat mechanics
- Monte Carlo
- molecular dynamics

## Favourite books

- "Understanding Molecular Simulation", Frenkel and Smit, Academic Press, 2002
- "Computer Simulation of Liquids", Allen & Tildesley, Oxford Science Publications, 1990

# Concepts

- potential energy
- work
- entropy
- free energy

## Goal

- There are models for potential energy
- for free energies we need simulations

# Nomenclature

- $T$  temperature
- $N$  all kinds of things, usually number of particles
- $k$  and  $k_B$  Boltzmann's constant
- $R$  gas constant =  $k N_a$
- $S$  entropy
- $F$  Helmholtz (NVT)
- $G$  (Gibbs) free energy (NPT)
- $U$  internal energy
- $H$  enthalpy =  $U + pV$
- $E$  energy,  $E_{kin}$ ,  $E_{pot}$ ,  $E_{bond}$ , ...
- $Q$  heat
- $W$  work
- $V$  volume
- $\Omega(x)$  number of states of system with property ( $x$ )
- $\vec{r}$  or  $\vec{r}$  vectors

# Potential energy

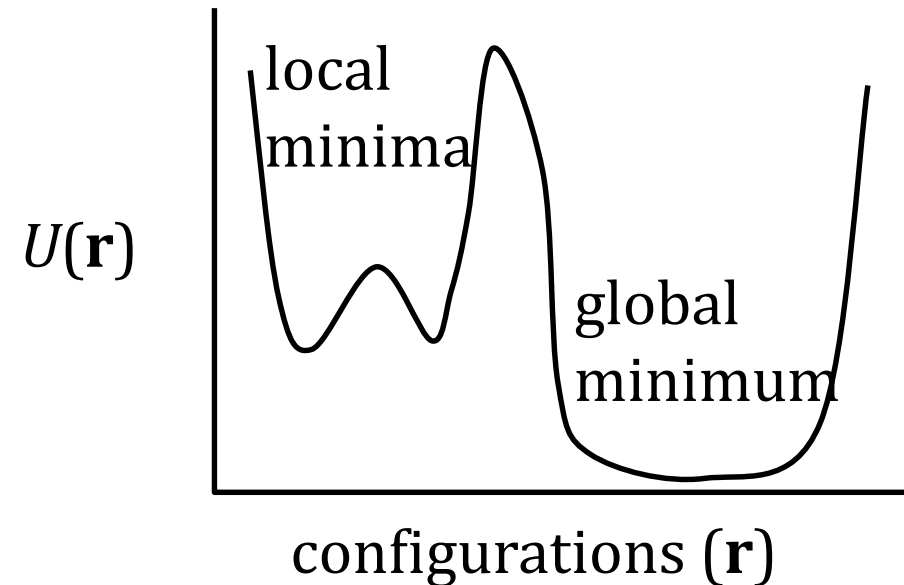
electrostatic  $U(r) = \frac{q_1 q_2}{4\pi\epsilon_0 r}$

gravity  $U(r) = \frac{Gm_1 m_2}{r}$

elastic, ...

potential energy of a protein...

- where would the protein like to be ?
- will be more interesting in ten minutes



# Conserving energy

For this course:

- first law
- relevant energies
- is it obvious ?
- when do we not conserve energy ?

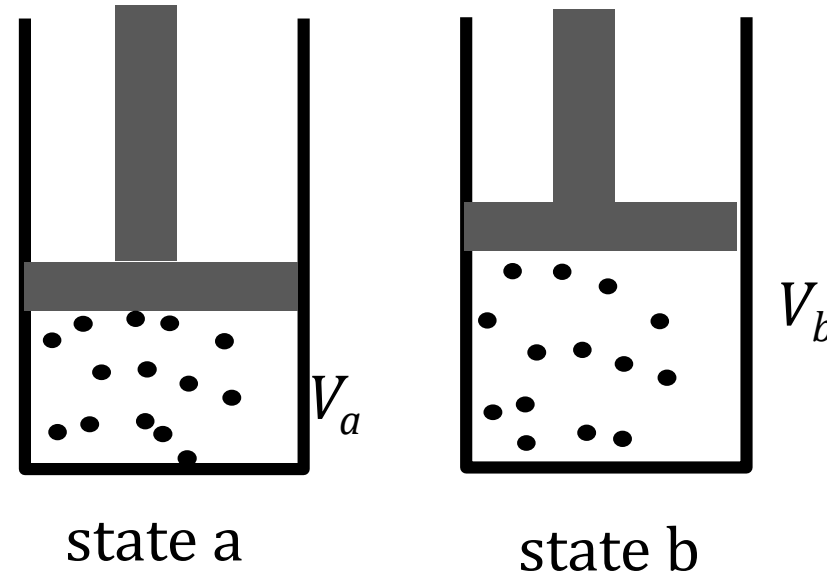
First law

- $dU = \Delta Q + \Delta W$
- $dU$  change in internal energy
- $\Delta Q$  heat given to system
- $\Delta W$  work done on system
- example of work...

# Work on a gas

$$\Delta W = -P\Delta V$$

$$\begin{aligned} dU &= \Delta Q + \Delta W \\ &= \Delta Q - P\Delta V \end{aligned}$$



## Others

- charges in electric fields, surface tension / area, ....
- we do not often do this in this course
- in one lecture:  $W$  will be important

# Conserving energy

Newtonian dynamics coming soon

- what is the energy of a particle ?

$$E_{tot} = E_{pot} + E_{kin}$$

- energy of a system ? (box of gas, protein in water)

$$E_{tot} = \sum_{i=1}^{N_{particles}} (E_{pot_i} + E_{kin_i})$$

- total energy is conserved – potential + kinetic

Do we always have kinetic energy ?

- Monte Carlo – no
- Wirkstoffentwurf-Vorlesungen - no

# Conserving energy – always obvious ?

- in the real world ? closed systems ? yes
- almost everything in this course – yes

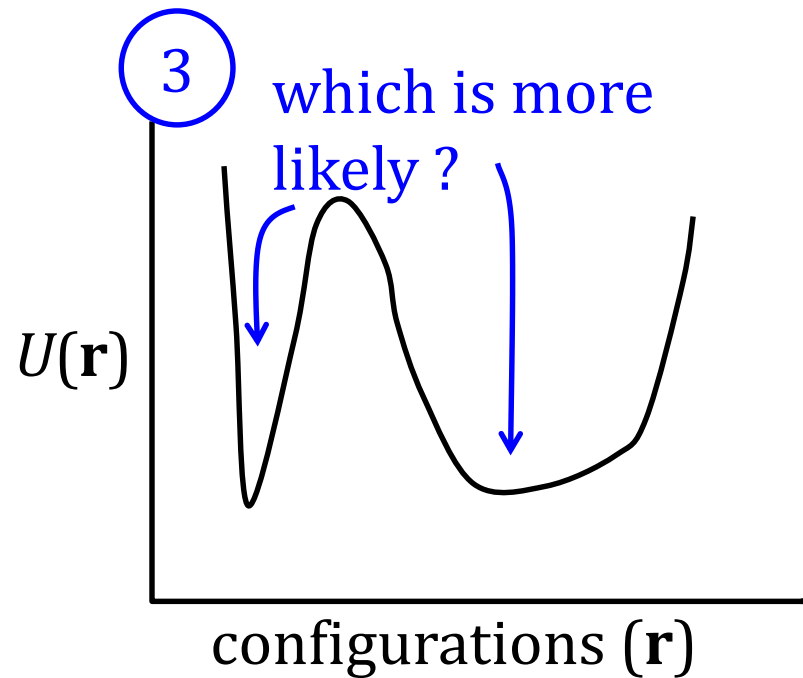
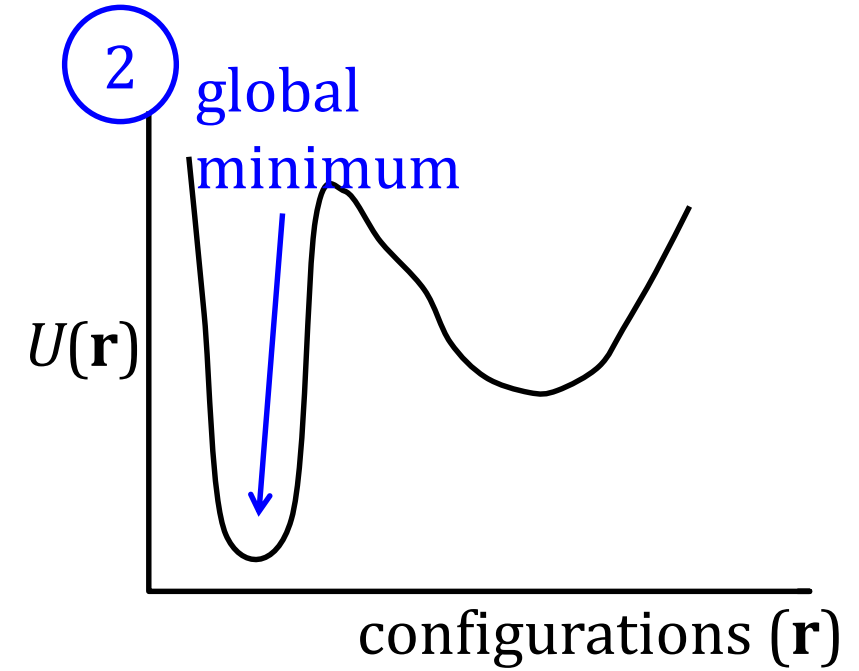
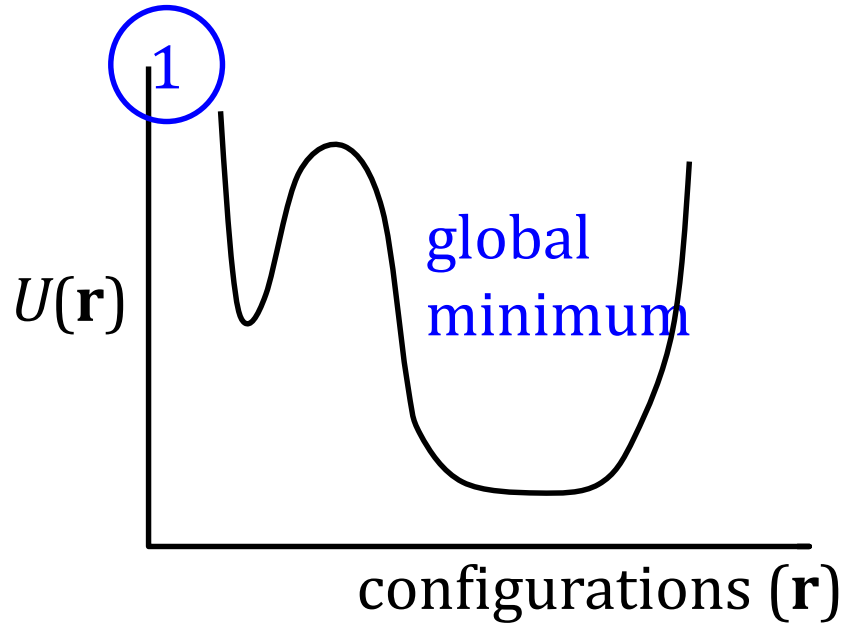
## Simulations

- if I want to use statistical mechanics, standard theory
  - no heating
  - no changing the force field (sounds obvious ?)
    - I want to simulate protein (un)folding
      - pull on the C-terminus – illegal – does work on system
      - heat system – illegal
    - push a drug (Medikament) into binding pocket – illegal

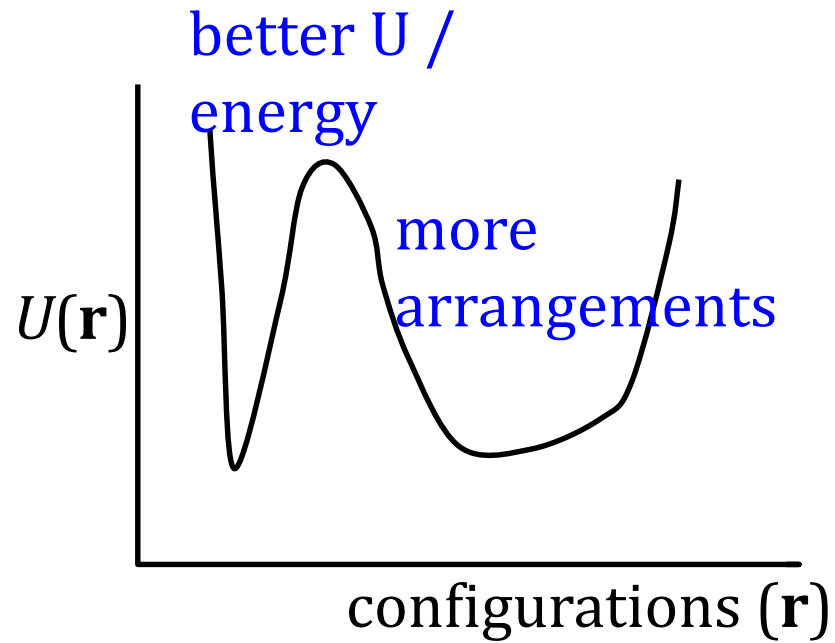
There is more than kinetic and potential energy..



# More than potential energy



# Simulation and formal statistical mechanics



Quantified ?

- entropy and free energy

Dumb simulation, preferring to go downhill

- should show how probabilities (entropy) are balanced against energy

# History

Statistical mechanics

- derived by summing up properties of individual particles

Thermodynamics

- less emphasis on individual particles

Lots of formulae which cannot be completely applied to proteins

- sums over infinite volumes, numbers of particles, time

## Rules and Limitations

Always at equilibrium

# Entropy

- you tend to become disordered (2<sup>nd</sup> law)
- role in free energy
- formula for simple systems

Disorder, how many ways can the system arrange itself..

- depends on energy (and volume and number of particles)
- two formulae coming... simple, more useful

How many states can the system have (specified  $E$ ) ?  $\Omega(E)$

Entropy

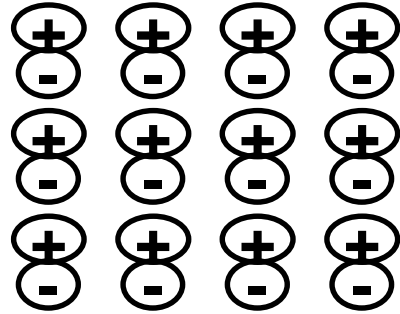
$$S = k \ln \Omega(E) \quad (\text{very simple form})$$

# Number of states $\Omega(X)$

$$T = 0$$

$$\Omega(E) = 1$$

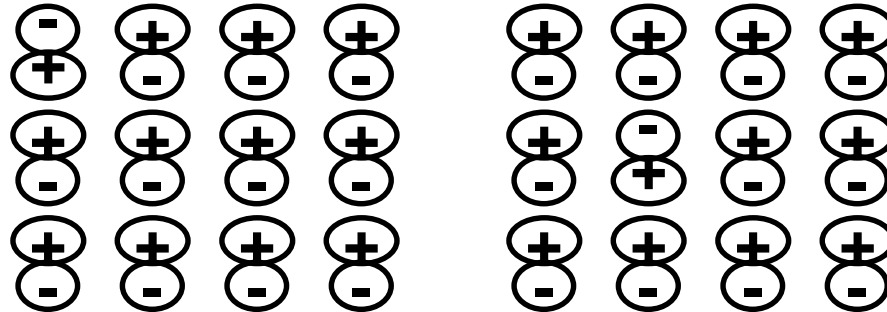
$$S = k \ln 1 = 0$$



$$T > 0 \quad \text{small}$$

$$\Omega(E) = n_{mol}$$

$$S = k \ln n_{mol}$$



A bit more energy ? more states

- more.. solid  $\rightarrow$  liquid .. many many more

# Gibbs

What if states are not equally likely ?

- $N_{state}$  states with distribution

Say  $p_1 = 0.999, p_2 = 0.00001, p_3 = 0.00001, \dots$

- just as if the system had one state - low entropy

Say  $p_1 = 0.5, p_2 = 0.1, p_3 = 0.1, \dots$

- a bit more entropy, but still very much dominated by  $p_1$

Say  $p_1 = 0.01, p_2 = 0.01, p_3 = 0.01, \dots$

- lots of states, all equally likely
  - lots of entropy

In general

$$S = -k \sum_{i=1}^{N_{state}} p_i \ln p_i$$

Gibbs entropy !!

# Applicability

Can one really estimate ?  $S = -k \sum_{i=1}^{N_{state}} p_i \ln p_i$

- liquid.. no
- protein ..no

With a simulation ?

- too many states
- approximations to  $\Delta S$  later

Simple system (grid / lattice)

- yes .. later in semester

# Units

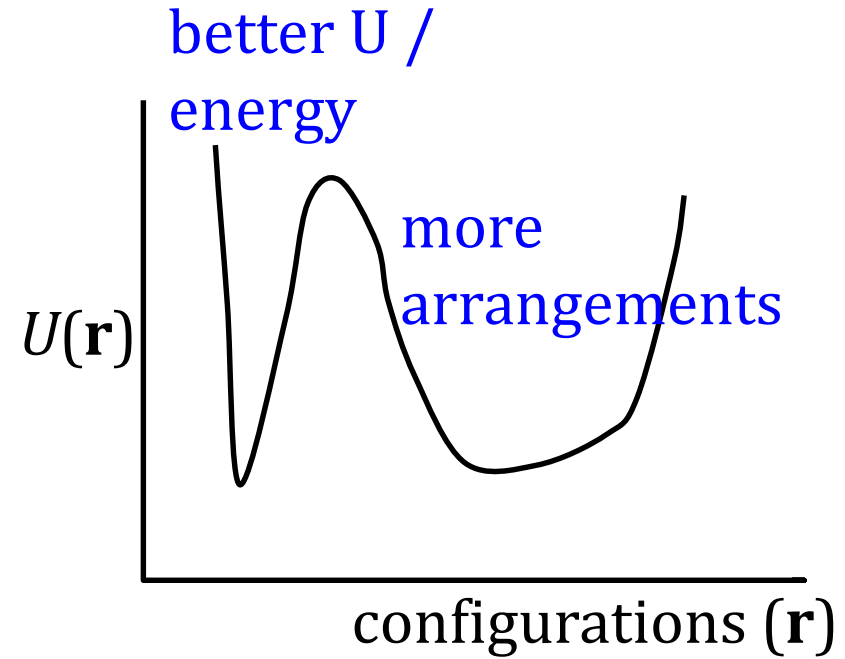
More arrangements = more entropy

Connected to energy

- balance of energy and entropy

Entropy units

- energy / temp ( $\text{JK}^{-1}$ )





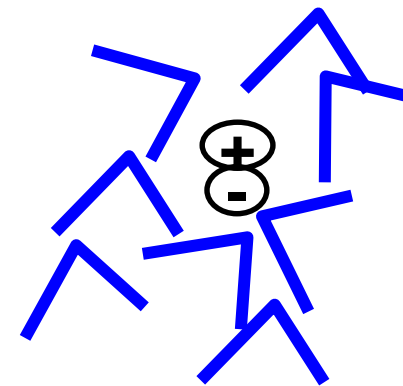
# Where can we use this ?

What is the entropy of a molecule ?

- sounds easy
- really...



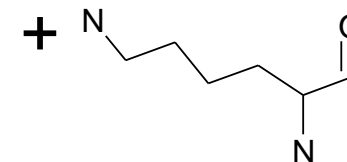
Entropy depends on solvent (system)



Entropy of a conformation ?

- has no meaning

$$S = -k \sum_{i=1}^{N_{state}} p_i \ln p_i$$



Is it additive ? Only rarely – not for interesting systems

# Adding entropy



$$\Omega_A$$



$$\Omega_B$$



$$\Omega_{AB} = \Omega_A \Omega_B$$

Entropy is "extensive"

$$S_{AB} = -k \ln(\Omega_A \Omega_B) = -(k \ln(\Omega_A) + k \ln(\Omega_B)) = S_A + S_B$$

Assumption

- for my new system A and B do not interact

What if they interact ?

- putting A in state 1 changes probability of B in state 1

$$S = -k \sum_{i=1}^{N_{state}} p_i \ln p_i$$

$$\Omega_{AB} \neq \Omega_A \Omega_B$$

# Adding entropy

I cannot usually add the entropy of two systems to get the entropy of new total system

Two systems

- could be boxes (previous slide)
- could be parts of a molecule (coming)

Do not add entropy terms ! Exam questions (promise)

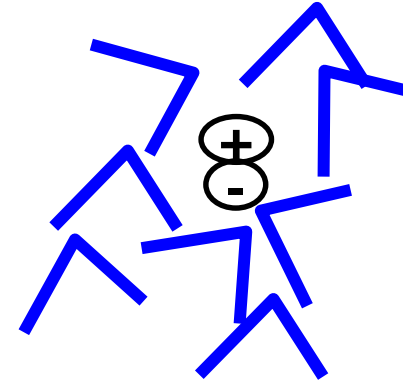
Can I decompose entropy ?

# May we decompose entropy ?

Remember energy in proteins

$$E_{tot} = E_{bonds} + E_{vdw} + E_{...}$$

$$S_{tot} = S_{bonds} + S_{vdw} + ... ?$$



No.

- Makes no sense unless bonds are decoupled from atoms and angles...

$$S_{ligand} + S_{solvent} + ...$$

- obviously do interact

Free energies

$$G_{solv}, G_{protein}, G_{ligand}, G_{bonds} \text{ not really legal}$$

- imagine bonds are system A, angles are system B, ..

# Decomposing Free Energy

Would you see this in the literature ?

Example

- in protein asp → asn mutation (COO<sup>-</sup> to CON)
- protein becomes more stable

Claim:

- the protein is more stable due to a change in electrostatic free energy

What has been implicitly assumed ?

There is no such thing as  $\Delta G_{el}$

# More examples – particle interactions

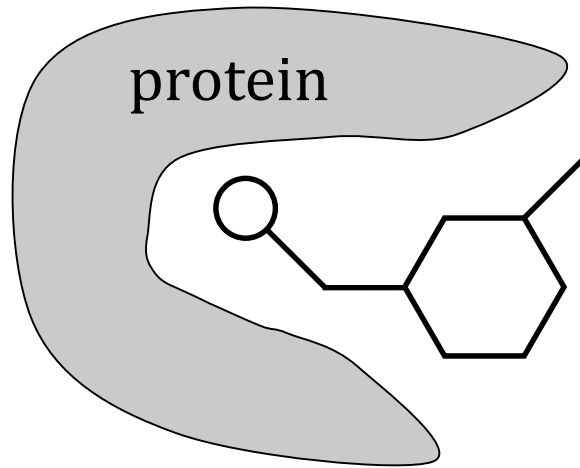
- You have a scoring function for interactions
- usually called  $\Delta G$

$$\Delta G_{tot} = \Delta G_{HB} + \Delta G_{ionic} + \dots$$

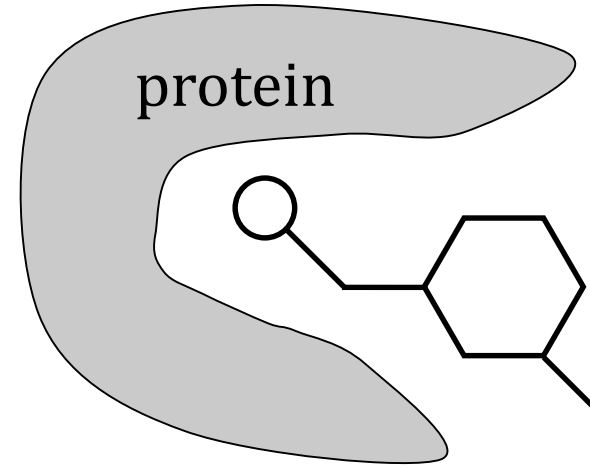
H-bonds  
ionic

- is this legal ?

# Entropy of a conformation ?



binding 1



binding 2

Can I talk about  $\Delta G_1$  vs  $\Delta G_2$  ?

$$S = -k \ln \Omega$$

but  $\Omega$  depends on all accessible states

If binding 1 can change to binding 2, they are both part of  $\Omega$

- I do believe we can talk about  $U_1$  and  $U_2$

# Meaning of labelled free energies

If we have "free energy contributions"

- we have entropy contributions

$\Delta G_{HB}$  implies  $H_{HB} + T\Delta S_{HB}$

$\Delta G_{ionic}$  implies  $H_{ionic} + T\Delta S_{ionic}$

what is  $S_{ionic}$ ?  $S = -k \ln \Omega_{ionic}$  ?

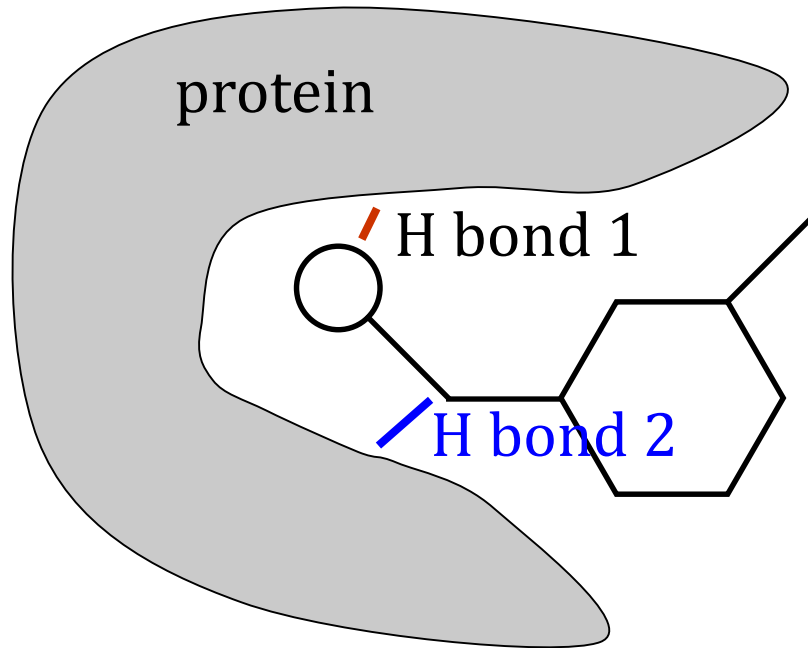
- no meaning

One cannot have "free energy contributions"



# independence of terms

Even without labelled entropies



Can we add free energy due to Hbond 1 and 2 ?

- implies  $\Delta S = \Delta S_1 + \Delta S_2$
- $-k \ln \Omega = -(k \ln \Omega_1 + k \ln \Omega_2)$
- $-k \ln \Omega = -k \ln(\Omega_1 \Omega_2)$

- only possible if there is no interaction between 1 and 2
- additivity of free energies is widely used (usually wrong)

# Stop and summarise

- Internal energy  $U$  – things like springs
- Enthalpy  $H$  – includes pressure
- Gibbs free energy  $G$  what we usually use
  - no discussion of different ensembles in this course
- Free energy lets us incorporate the concept of what is most likely
- Entropy is not additive between systems (parts of systems) that interact

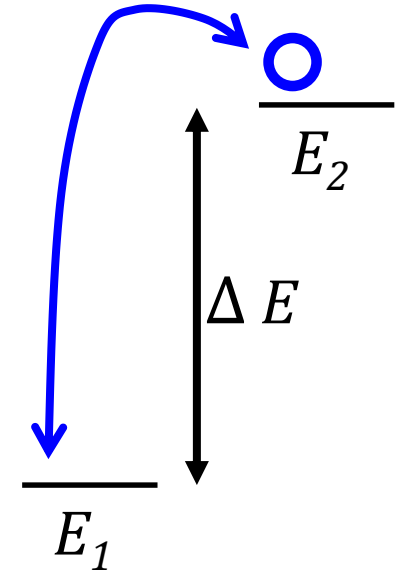
# Boltzmann distribution – two states

System with two energy levels

- I know the temperature and energy difference
- How likely is system to sit in  $E_1$  ?  $E_2$  ?
  - System is very cold,
    - $\Delta E$  seems big
  - System is very very hot  $T \rightarrow \infty$ 
    - $\Delta E$  does not matter, seems small
- Relative populations (probabilities)  $p_1, p_2$
- $\Delta E = E_2 - E_1$

$$\frac{p_1}{p_2} = e^{-\Delta E / kT}$$

- sometimes  $\frac{p_1}{p_2} = e^{-\beta \Delta E}$



# Sign conventions

$$\Delta E = E_1 - E_2 \quad \text{or} \quad \Delta E = E_2 - E_1 \quad ?$$

For exams

- the lower energy state is always more populated
- there are always questions about probabilities / energies
- state with most negative energy is most populated

# Boltzmann distribution

What is the probability of a certain energy level ?

- depends on all available levels

$$p_i = \frac{e^{-\beta E_i}}{\sum_j e^{-\beta E_j}}$$

- name of bottom ... partition function,  $Z = \sum_i e^{-E_i/kT}$
- does this agree with previous slide ?

# Going from two to $n$ states

Start with

$$= e^{-\beta \Delta E}$$

$$\begin{aligned} \frac{p_1}{p_2} &= e^{-\beta (E_1 - E_2)} \\ &= \frac{e^{-\beta E_1}}{e^{-\beta E_2}} \end{aligned}$$

but this is for the ratio of two states. If we have  $n$  states, we are interested in

$$\frac{p_1}{p_1 + p_2 + p_3}$$

Let us turn this upside down and start with  $\frac{p_1 + p_2 + p_3 + \dots}{p_1}$

let us consider some more states and get the absolute fraction for  $p_1$

$$\begin{aligned} &= \frac{p_1}{p_1} + \frac{p_2}{p_1} + \frac{p_3}{p_1} + \dots \\ &= \frac{e^{-\beta E_1}}{e^{-\beta E_1}} + \frac{e^{-\beta E_2}}{e^{-\beta E_1}} + \frac{e^{-\beta E_3}}{e^{-\beta E_1}} + \dots \\ \frac{p_1 + p_2 + p_3 + \dots}{p_1} &= \frac{e^{-\beta E_1} + e^{-\beta E_2} + e^{-\beta E_3} + \dots}{e^{-\beta E_1}} \\ &= \frac{\sum_j e^{-\beta E_j}}{e^{-\beta E_1}} \end{aligned}$$

turn upside-down,

$$p_i = \frac{e^{-\beta E_i}}{\sum_j e^{-\beta E_j}}$$

where the summation runs over all states of the system



# Consequences of Boltzmann distribution

At absolute zero

- only lowest energy state is populated

At low temperatures

- low energy states favoured

High temperature

- system can visit high energy regions

Infinite temperature

- all states equally likely

For two states (bound / unbound)

- exponential term means populations quickly become big/small

$$\frac{p_1}{p_2} = e^{-\Delta E/kT}$$

# How big are energy differences ?

Simplest case, equal populations  $\frac{p_1}{p_2} = e^{\frac{\Delta E}{kT}}$

$$p_1 = p_2$$

$$1 = e^{\Delta E/kT}$$

$$\ln 1 = \Delta E/kT$$

$$\Delta E = 0$$

$k = 1.38 \times 10^{-23} \text{ J K}^{-1}$ , but use

$$\begin{aligned} R &= k N_A \\ &= k \cdot 6.02 \times 10^{23} \text{ J K}^{-1} \text{ mol}^{-1} \\ &= 8.314 \text{ J K}^{-1} \text{ mol}^{-1} \end{aligned}$$

# Examples of populations

$$\frac{p_1}{p_2} = e^{\frac{\Delta E}{kT}}$$

For 99:1 at 300 K

- $\Delta E = 11 \text{ kJ mol}^{-1}$

$$\ln \frac{p_1}{p_2} = \frac{\Delta E}{kT}$$

$$\Delta E = kT \ln \frac{p_1}{p_2}$$

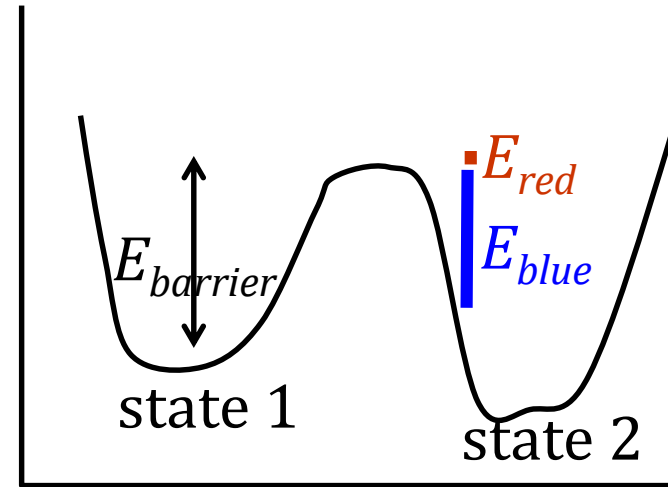
## Drugs

- Dissociation constant  $k_d$  of drug D to protein P  $k_d = \frac{[D][P]}{[DP]}$
- nanomolar drugs are  $10^{-9}$ ,
- difference between 1 and 10 nm binding drug
  - 46 vs 52  $\text{kJ mol}^{-1}$
- topic will return later

# Barrier crossing

How likely are you to cross a barrier ?

- $p_{red}$  vs  $p_{blue}$



- all the blue copies of system will not make it over  $E_{barrier}$
- red population is small, even for  $E_{barrier} \sim kT$
- but explains why  $\ln(rate) \propto T$

## **where next ?**

- how to see some of these properties by simulating
- how simulating depends on these properties