

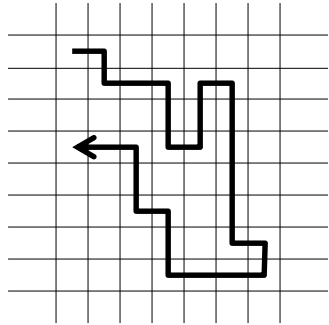
# Monte Carlo and MD simulations

Andrew Torda, April 2018 strukt und sim

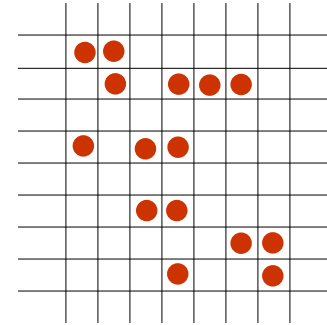
What we observe in any system ?

- averages of observables (pressure, energy, density)

Given enough time system will visit all states



time



random  
hopping

My observable  $\mathcal{A}$

$$\mathcal{A}_{obs} = \frac{1}{b-a} \int_a^b \mathcal{A}_t dt$$

$$\mathcal{A}_{obs} = \frac{1}{N_{obs}} \sum_{i=1}^{N_{obs}} \mathcal{A}_i$$

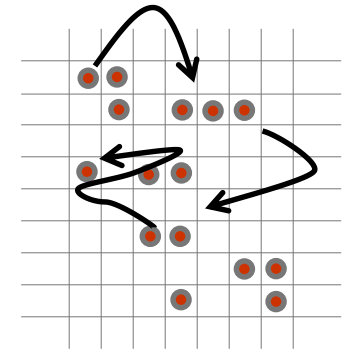
# Time and space averages

If we believe  $\mathcal{A}_{obs} = \frac{1}{N_{obs}} \sum_{i=1}^{N_{obs}} \mathcal{A}_i$

then

$$\begin{aligned} \mathcal{A}_{obs} &= \sum_j^{states} p_j \mathcal{A}_j \\ &\equiv \langle \mathcal{A} \rangle \end{aligned}$$

and  $p_j$  is the probability of state  $j$



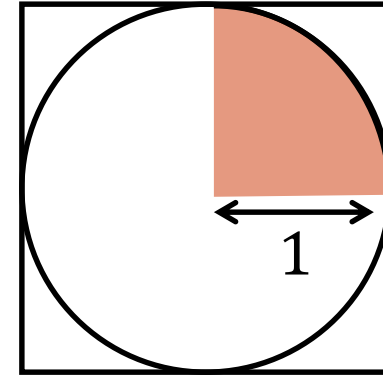
- $\langle \mathcal{A} \rangle$  is ensemble average and usually  $\bar{\mathcal{A}}$  is time average
- if sample with correct probability, we can find  $\mathcal{A}_{obs}$
- order of visiting states does not matter

# Monte Carlo

How to calculate  $\pi$  with random numbers

$$\frac{\text{points}_{red}}{\text{points}_{square}} = \frac{1/4 \pi r^2}{\text{area in square}}$$

$$\pi = 4 \frac{\text{points}_{red}}{\text{points}_{square}}$$



```
while ( not converged)
  pick random x, y
  n_square++
  if ((x2+y2) < 1)
    n_red++
print  $\frac{4 n_{red}}{n_{square}}$ 
```

# Generating distributions / Monte Carlo

Generating points in a circle ? (generating function)

$$p_{in\_circle} = \begin{cases} 1 & x^2 + y^2 \leq 1 \\ 0 & x^2 + y^2 > 1 \end{cases}$$

We could work out the area of a circle (integrate) by picking random numbers

What does Monte Carlo simulation mean ?

- generating points according to some distribution to find an average or integral
- what is our distribution in physical systems ?
  - Boltzmann distribution

# Monte Carlo and Boltzmann distributions

Boltzmann probability distribution

$$p_i = \frac{e^{\frac{-E_i}{kT}}}{\sum_j e^{\frac{-E_j}{kT}}} \quad \text{often written as} \quad p_i = \frac{e^{\frac{-E_i}{kT}}}{Z} \quad \text{since we define } Z = \sum_j e^{\frac{-E_j}{kT}}$$

If we could generate this distribution,  
we could reproduce most properties of a system

Leads to a scheme (not possible)

## correct, but not practical scheme

while (not happy)

    generate configuration  $\mathbf{r}_i$  (conformation of protein, ...)

    calculate  $p_i$  (number between 0 and 1)

    generate random number  $x$

    if ( $x < p_i$ )

        accept  $\mathbf{r}_i$

    else

        reject  $\mathbf{r}_i$

$$p_i = \frac{e^{\frac{-E_i}{kT}}}{\sum_j e^{\frac{-E_j}{kT}}}$$

- Result ? a set of  $\mathbf{r}_i$  with Boltzmann distribution
- Problem ? we do not know  $\sum_j e^{\frac{-E_j}{kT}}$

## a better scheme

We cannot generate points from  $p_i = \frac{e^{-E_i/kT}}{\sum_j e^{-E_j/kT}}$

What if we have two configurations ?

$$\frac{p_i}{p_j} = \frac{e^{-E_i/kT}}{Z} \frac{Z}{e^{-E_j/kT}}$$

$$= e^{\frac{E_j - E_i}{kT}}$$

$$= e^{\frac{-\Delta E}{kT}}$$

# a better scheme

$$\frac{p_i}{p_j} = e^{\frac{-\Delta E}{kT}}$$

If we have one configuration to start

- we can work out the relative probability of a second

Convenient convention

- going from old  $\rightarrow$  new  $\Delta E < 0$ 
  - $E_{new} - E_{old} < 0$  energy is better / more negative

Does it matter where you start? What is  $i$ ?



# Metropolis Monte Carlo

- generating a distribution  $\frac{p_i}{p_j} = e^{\frac{-\Delta E}{kT}}$
- if  $\Delta E < 0$ , new is likely (more than 1)
- if  $\Delta E > 0$ , old is  $p_{new}$  is possible

```
generate starting configuration  $\mathbf{r}_o$ 
while (not happy)
    generate  $\mathbf{r}_{new}$ 
    calculate  $E_{new}$  and  $\Delta E$ 
    if  $\Delta E < 0$ 
        set  $\mathbf{r}_o$  to  $\mathbf{r}_{new}$ 
    else
        x = rand [0:1]
        if( $x \leq e^{-\Delta E/kT}$ )
            set  $\mathbf{r}_o$  to  $\mathbf{r}_{new}$ 
```

- what if  $\Delta E$  slightly  $> 0$  ?
  - 0.0000000001
- what if  $\Delta E = 10^6$  ?
- small uphill moves are OK
- bigger moves are less likely

# Properties of Monte Carlo

The set of  $\mathbf{r}_o$  is a valid distribution (ensemble)

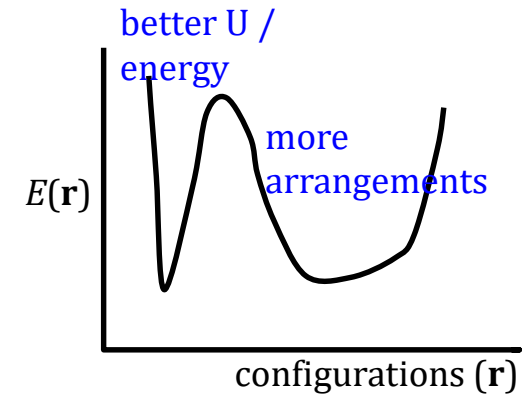
- for some property  $\mathcal{A}$

$$\mathcal{A}_{obs} = \langle \mathcal{A} \rangle = \frac{1}{N_{visited}} \sum_i^{N_{visited}} \mathcal{A}_i$$

- $\mathcal{A}$  could be density, structural property,  $E$ , ...
- only works for one temperature  $T$

Look at picture.. could I calculate entropy / free energy ?

- for simple systems



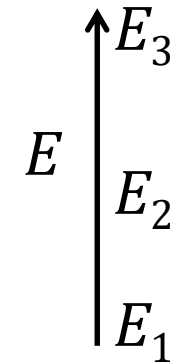
# Equilibrium

MC results (observables / averages)

- only for system at equilibrium
- simulations generate system at equilibrium

What happens for a system out of equilibrium ?

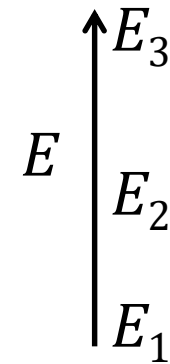
- Toy system with 3 states
- for some  $T$ , at equilibrium
- $p_1 = 5/8$        $p_2 = 1/4$        $p_3 = 1/8$
- if I have 80 copies of the system, most are in state<sub>1</sub>



# Reaching equilibrium

System wants  $p_1 = 5/8$      $p_2 = 1/4$      $p_3 = 1/8$   
50 : 20 : 10

- start it with 5 : 70 : 5
- all moves  $2 \rightarrow 1$  are accepted (large flux)
- the flux from  $1 \rightarrow 2$ 
  - $1 \rightarrow 2$  moves are not always accepted
  - there are less particles in state<sub>1</sub>



Moving to equilibrium depends on

- population
- probability

# Detailed balance

For any two states (state<sub>*i*</sub> and state<sub>*j*</sub>)

Flow  $i \rightarrow j$  must equal  $j \rightarrow i$

- otherwise ?

Flow  $i \rightarrow j$  depends on

- population  $N_i$
- probability  $\pi(i \rightarrow j)$

Detailed balance

$$N_i \pi(i \rightarrow j) = N_j \pi(j \rightarrow i)$$

- detailed balance must apply for any pair  $i, j$

all textbooks use  $\pi$  for probability here

# Ergodic

## Assumptions

- I can do integrals because
  - I will visit every state
  - I can calculate  $p_i$  for all states
- I will visit every state

alternatively

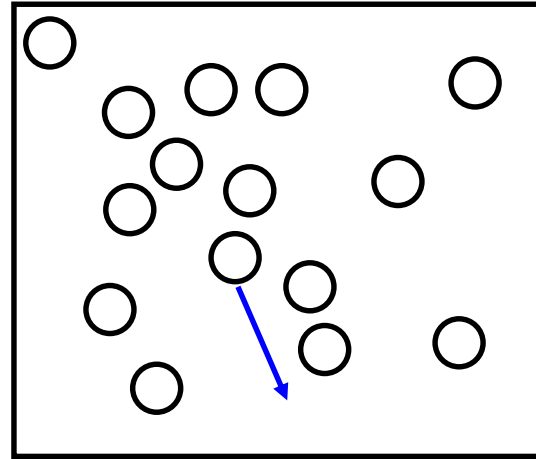
For any  $i, j$

- $\pi(i \rightarrow j) > 0$
- may require a finite number of steps:  $i \rightarrow k \rightarrow m \rightarrow j$
- must be satisfied

# Moves

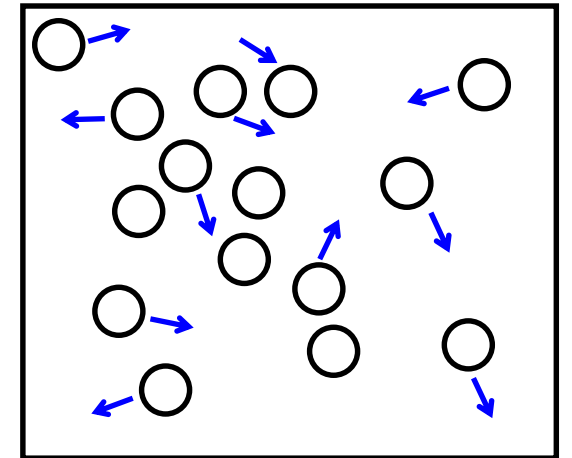
## version 1

- decide on  $r_{max}$
- pick a particle at random
- pick random  $\Delta x, \Delta y, \Delta z$   
 $0 < \Delta a < r_{max}$
- apply move
- accept / reject move



## version 2

- decide on smaller  $r_{max}$
- foreach particle
  - pick random  $\Delta x, \Delta y, \Delta z$   
 $0 < \Delta a < r_{max}$
- apply move
- accept / reject



# Moves

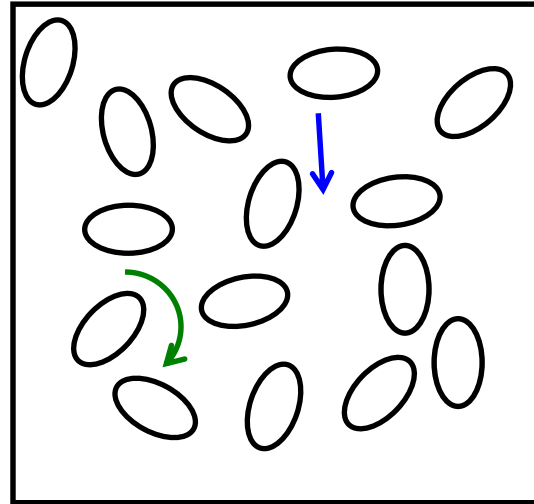
- both kinds of move OK
- note
  - "accept / reject"

More generally,

- how big is  $r_{max}$  ?
- big
  - system moves faster
  - more moves rejected

What if my particles are not spheres ?

- rotations also necessary
- time has no meaning





# Bonded systems

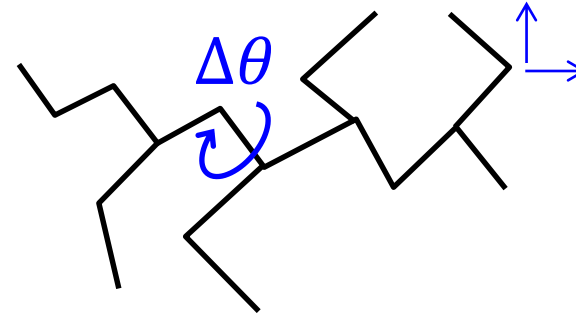
Protein (lipid, polymer, ..)

Random  $\Delta x$  ?

- nearly all will stretch a bond
  - high energy : rejected move
- only feasible method
  - random rotations  $\Delta\theta$

In general

- most kinds of simple moves OK
- must maintain detailed balance, ergodicity
- question of efficiency
  - high rejection rate means lots of wasted calculations



# More moves - $N$ particles

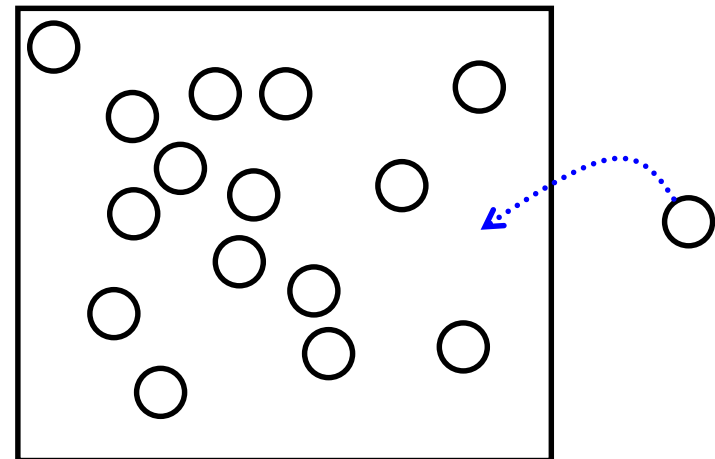
$$\frac{p_{new}}{p_{old}} = e^{-\Delta E/kT}$$

I have defined temperature

- and  $N_{particles}$  and  $V$
- called NVT simulation

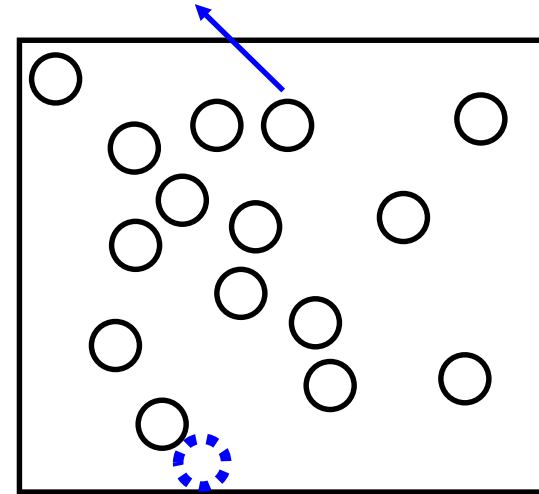
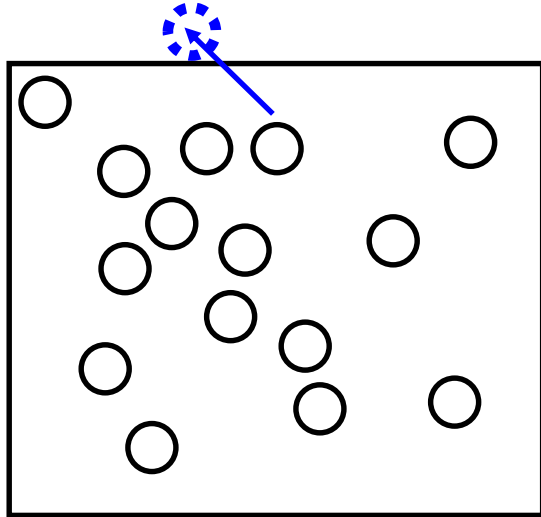
Could I have varied something else ?

- what if I tried to put particles in / take out ?
  - sometimes energy  $\uparrow$  sometimes  $\downarrow$
- system will fluctuate around  $\langle N \rangle$
- this would not be NVT

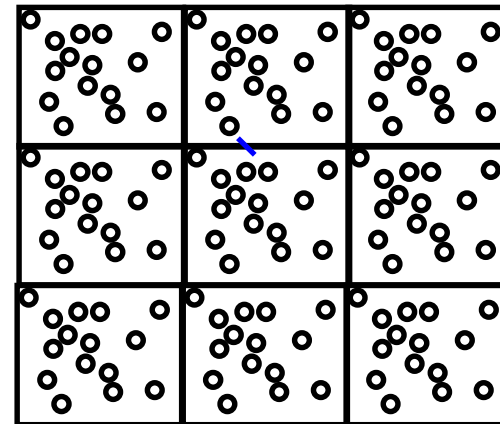


# periodic boundary conditions

Relevant to gases, proteins in water, ..



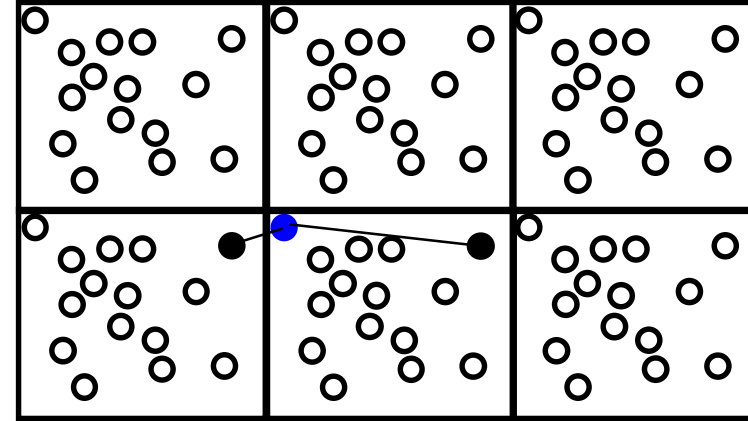
Behaves like an infinite system



# Infinite interactions ?

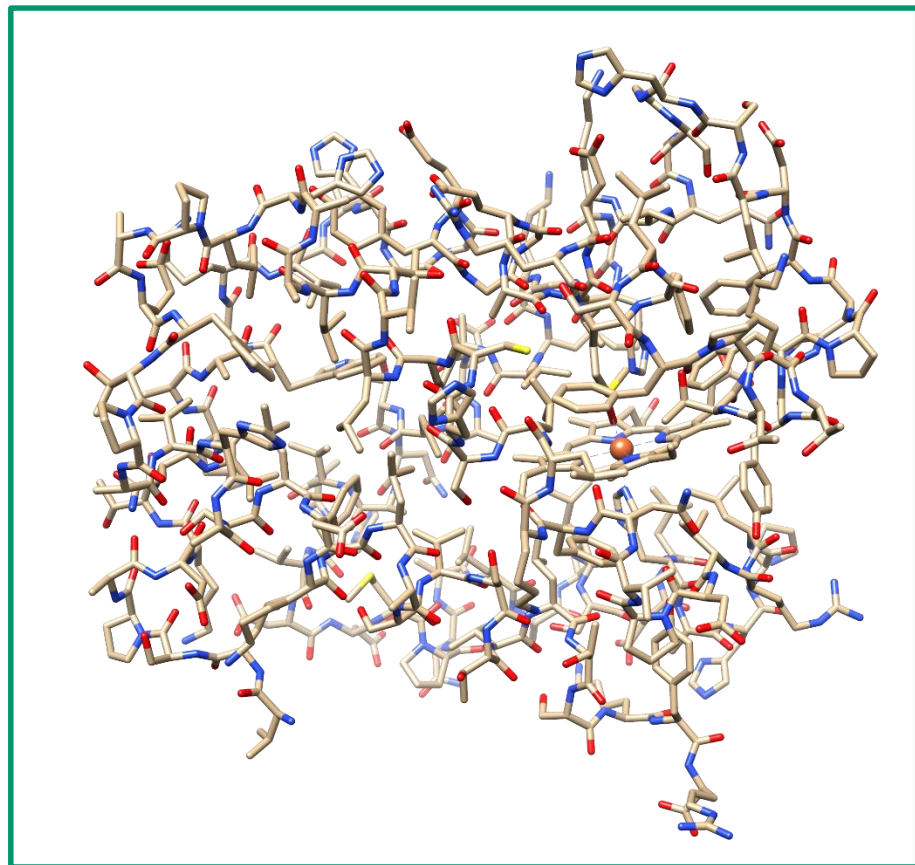
Neighbours of blue particle

- only use the nearer
- not really an infinite system
- volume defined by box




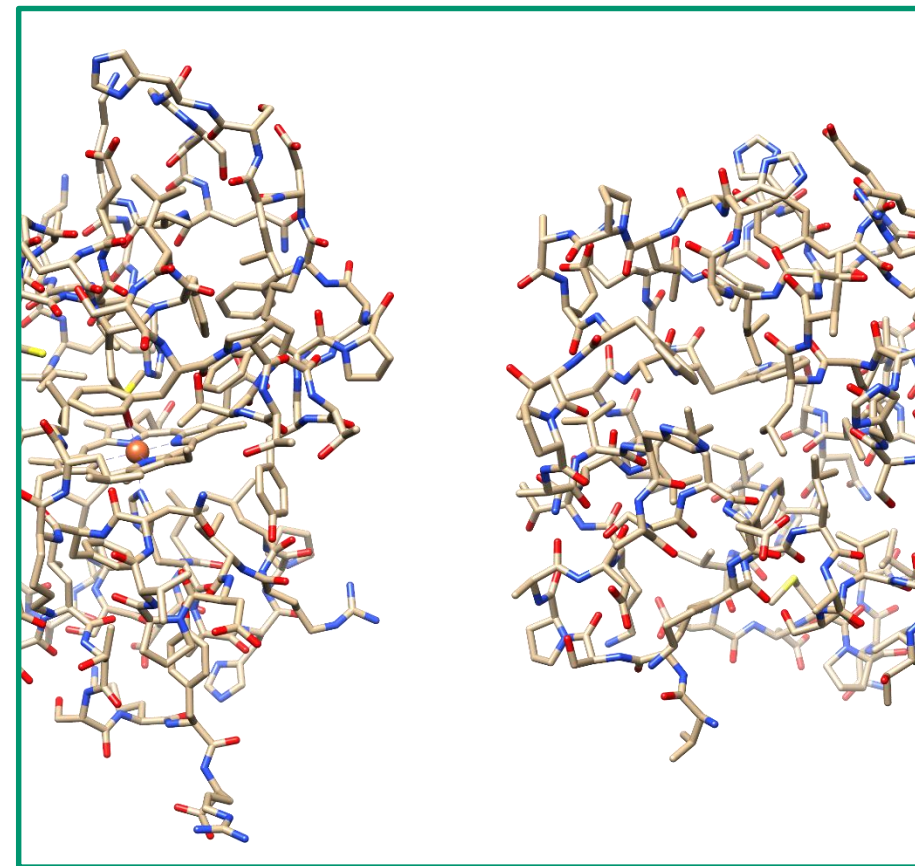
- how many neighbours does blue particle have ?
- why do we need cutoffs ?

# protein chopped up by periodic boundary conditions



start  
coordinates

simulation  
diffusion

A green arrow pointing from the 'start coordinates' box to the 'later' box, indicating the progression of the simulation.

later

# Problems with Monte Carlo

while (not happy)  
    propose move  
    accept / reject move

Small steps ?

- system moves slowly: long time to visit all states

Big steps ?

- calculate energy
- reject move
  - no progress, wastes time

# Dense Systems and Monte Carlo

Random moves ?

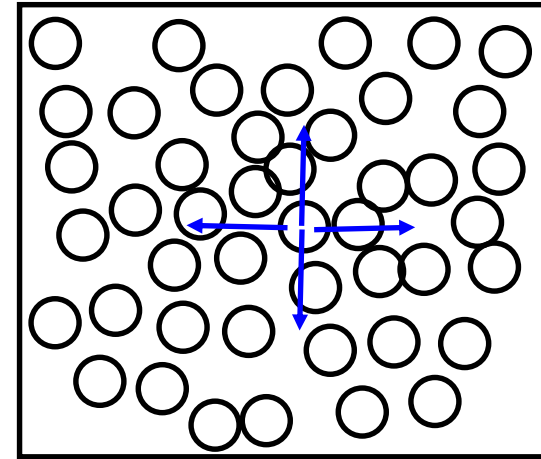
- most moves rejected

Dense systems ?

- liquids
- proteins, polymers, ...

Solutions

- cleverer MC moves (later)
- MD



# Why do molecular dynamics simulations ?

## Real world

- box of gas, molecule in space, protein molecule in water
- atoms hit each other,
  - share energy, box expands/contracts, ..
  - soon reaches equilibrium
  - visits low energies (often), high energies (less often)
  - visits entropically favoured regions
- we stick in a thermometer
- measure density, ...

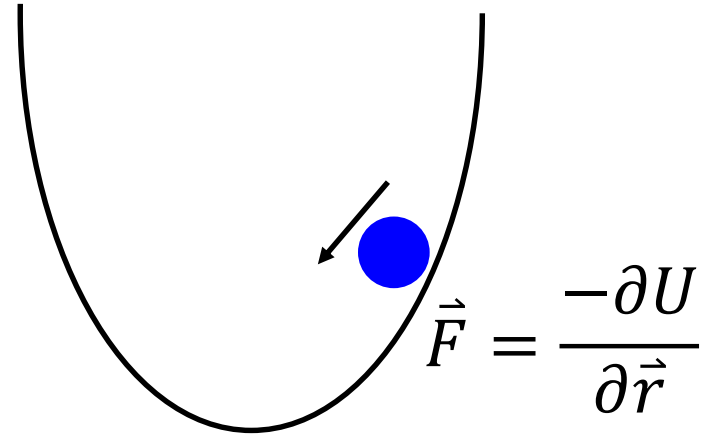
## What have the atoms done ?

- feel forces and move
- an MD simulation just copies this



# What do we expect ? Molecular Dynamics

one particle in a well



Unlike MC, particles have kinetic energy  $E_{kin}$

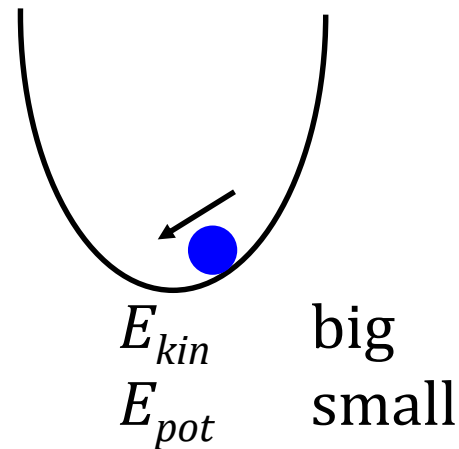
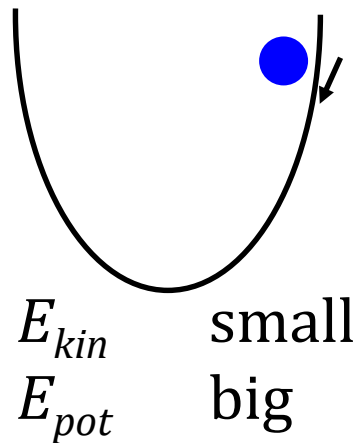
# Kinetic and potential energy

Our system is isolated (no work done)

$E_{tot}$  never changes

- conserves energy (no work done on system)

$$E_{tot} = E_{pot} + E_{kin}$$



For one particle  $E_{tot} = E_{pot} + E_{kin} = \text{constant}$

# Lots of particles

Particles hitting each other

- exchanging energy

Total system

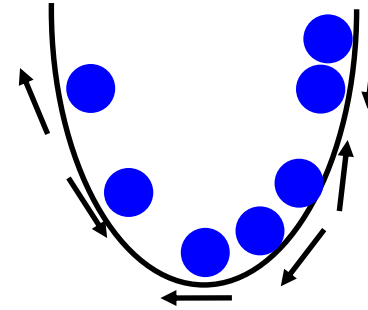
- conserves energy

One particle ?

- maybe at bottom but moving slow ( $E_{kin} + E_{pot}$  small)
- per particle energy no longer conserved (may gain or lose)

Many particles

- distribution of velocities
- distribution of potential energies



# Boltzmann distribution in real world

One version of real world (N, V, T)

- constant number of particles, volume, temperature
- today  $E = E_{kin} + E_{pot}$
- $Z$  is partition function
- earlier  $Z = \sum_i e^{\frac{-\Delta E_i}{kT}}$

But now we have kinetic energy  $E_{kin}(\mathbf{p})$

- where  $\mathbf{p} = m\dot{\mathbf{x}}$ 
  - potential energy  $E_{pot}(\mathbf{r})$
- if we write in continuous form ...

# Partition function for MD

Usually write  $\mathcal{H}(\mathbf{p}, \mathbf{r}) = E_{kin}(\mathbf{p}) + E_{pot}(\mathbf{r})$

- "Hamiltonian"

All the states are defined by all possible momenta and coordinates

- sum over these:  $Z(N, V, T) \propto \int d\mathbf{p} \int d\mathbf{r} e^{\frac{-\mathcal{H}(p,r)}{kT}}$

often see  $H(\mathbf{p}, \mathbf{r})$  or  $\mathcal{H}(\Gamma)$

# MD Method

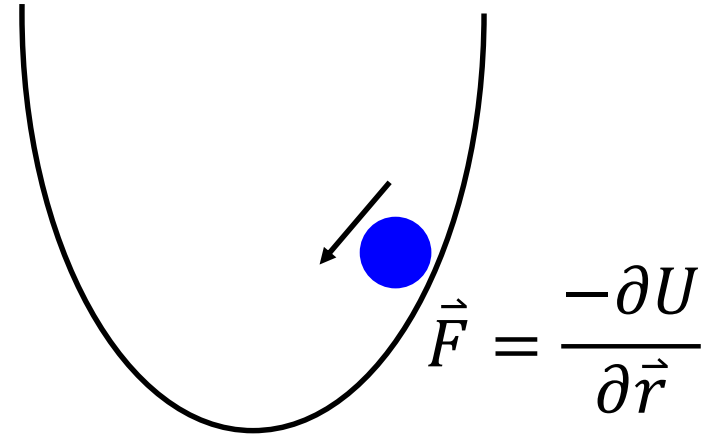
For any particle we can calculate forces

Newtons law

$$F = ma \text{ often better written } \ddot{\vec{x}} = \vec{F}m^{-1}$$

If we know acceleration

- we can get velocity
- from velocity
- can get coordinates



```
while (nstep < max_step)
    calculate forces
    integrate to get new coordinates
    nstep ++
} averaging,
sampling,
...
```

# starting a system

## Initial coordinates

- protein model
- protein from protein data bank (PDB)
- protein + proposed ligand
- box of liquid

## Do initial coordinates matter ?

- in principle: no
  - infinately long simulation visits all configurations, reaches equilibrium
- in practice: yes
  - bad examples
    - no simulation is long enough to predict protein conformation
  - take water configuration and run at ice temperature

# Initial velocities

First consider temperature – reflects kinetic energy

$$\left\langle \frac{1}{2} m v_{\alpha}^2 \right\rangle = \frac{1}{2} kT$$

where  $v_{\alpha}^2$  could be  $v_x, v_y, v_z$

leads to definition

$$T(t) = \sum_{i=1}^N \frac{m_i v_i^2(t)}{k N_f}$$

- where  $N_f$  is number degrees of freedom  $\approx 3N$
- we could use this to get initial velocities  $\langle v_{\alpha}^2 \rangle = \frac{kT}{m}$



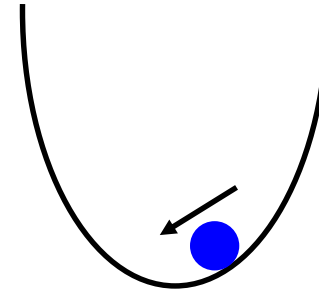
# Initial velocities

Would one  $\langle v^2 \rangle$  be OK ?

- not very good
  - $E_{kin}$  correlated with  $E_{pot}$

Either

- use more sophisticated distribution
- do not worry
  - system will go to equilibrium
    - velocities will reach sensible values



# Getting new velocities / coordinates

constant acceleration

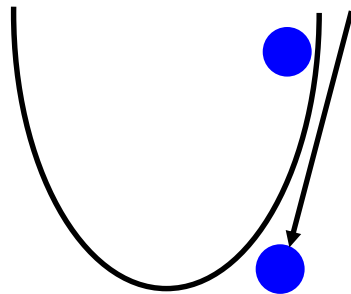
$$x_t = x_0 + vt + \frac{1}{2}at^2$$

or

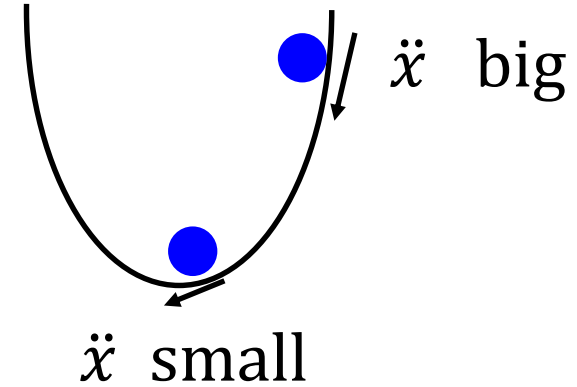
$$x_t = x_0 + \dot{x}t + \frac{1}{2}\ddot{x}t^2$$

OK for constant acceleration

- try to use formula to predict future time



big  $\Delta t$  / step  
big error



small  $\Delta t$  / step  
small error  
slow

# Fundamental problem with integration

- We want to use big  $\Delta t$  (speed)
- We must use small  $\Delta t$  (accuracy)

All  $\Delta t$  will give us some error

- numerical integration is never perfect

How small is  $\Delta t$  ?

- depends on fastest frequency / steepest walls in energy
  - usually bonds
- for proteins at room temperature
  - $\Delta t \approx 1$  fs (femtosecond  $10^{-15}$  s)
- high temperature  $\Delta t$  should be smaller

# Noise and heating

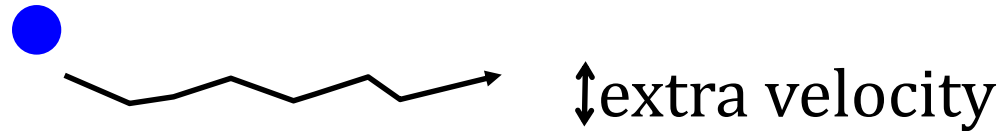
## General rule

- noise heats the system
- formally difficult to prove
- $E_{kin} = \frac{1}{2}mv^2$

● no kinetic energy



●↕  $E_{kin}$  due to noise



# Noise-free Simulation

Energy conservation : Absolute rule  $E_{pot} = f(\mathbf{r})$

- no time component
- invariant under translation, rotation

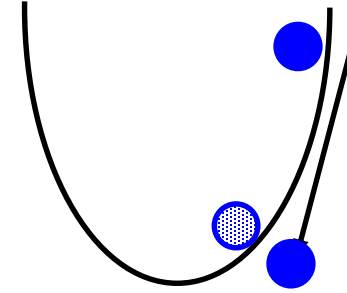
When violated ?

- $(\mathbf{r})$  does not change, but  $E_{pot}$  changes:  $E_{tot}$  changes

# Noise Sources

## Integrator

- coordinates do not match velocity  
 $E_{kin}$  wrong:  $(E_{kin} + E_{pot}) \neq \text{constant}$
- energy not conserved

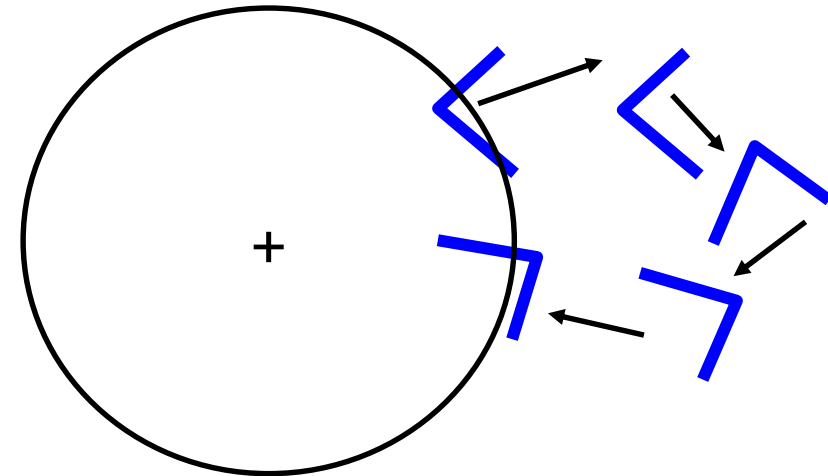


## Numerical noise

- $E_{pot} = f(\mathbf{r})$
- initial coordinates ( $\mathbf{r}$ ) quoted to 3 decimal places

## Cutoffs

- within cutoff rotation restricted
- outside cutoff rotation suddenly free



## Result

- heating

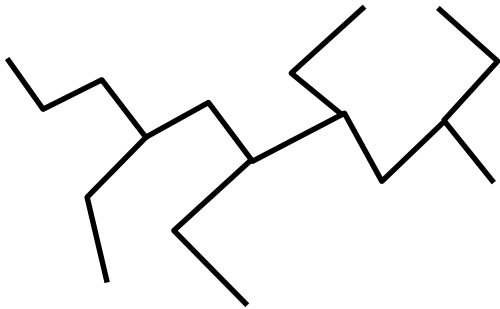
# Equilibrium

Remember MC story

- system not at equilibrium ? eventually equilibrates

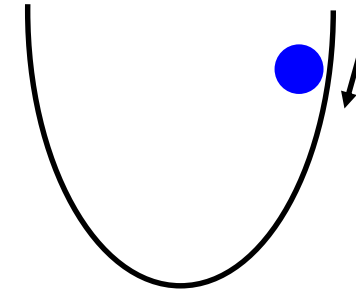
MD

- start in high energy  $E_{pot}$
- $E_{pot}$  converted to  $E_{kin}$



Some high energy conformation

- relaxes
- $E_{pot}$  converted to  $E_{kin}$



MD system will not

- really find low energy
- known temperature

# MD in a closed system

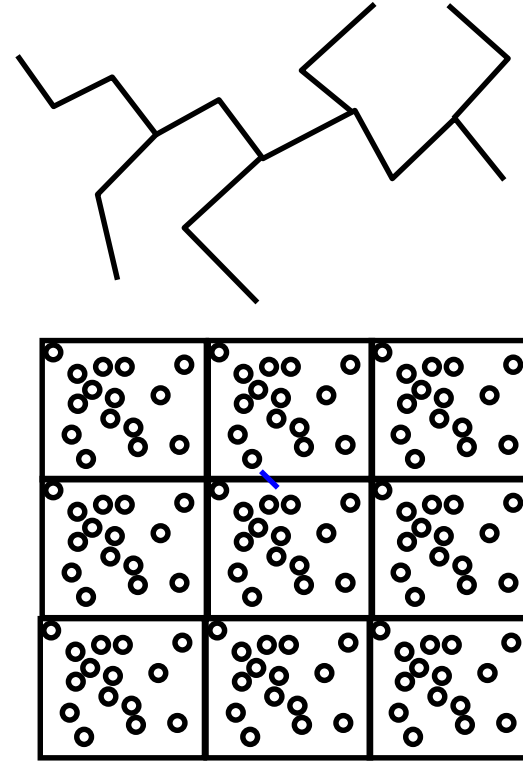
- An isolated molecule should not lose energy
- A repeated box will not lose energy
- Formally system is
  - NVE (constant  $N_{particles}$ , volume, energy)

## Problems

- we want to set the temperature of the system
- we may have noise / heat creating energy

## Cure

- thermostat

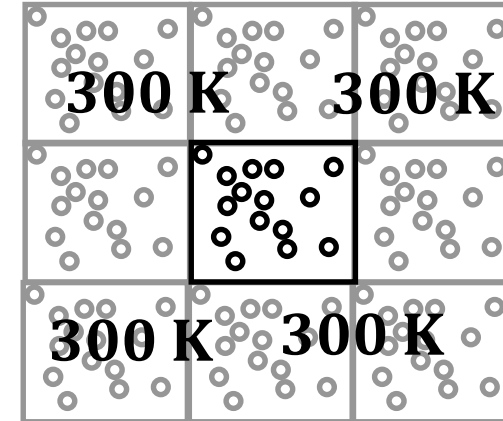




# Bath

Imagine infinite bath at desired temperature

- heat will flow in or out
- at equilibrium no flow of heat
  - maybe removal of noise/heat



How to implement? Many ways

Occasionally:

1. introduce a fake particle desired temperature / collide
2. pick a particle at random / give average  $v$  for temperature
3. Easy method –weak coupling...

# Weak Coupling

Remember temperature\*  $E_{kin} = \sum_i^N \frac{1}{2} m_i v_i^2 = \frac{3}{2} NkT$

Goal: heat leaves system depending on how wrong temperature is

$$\frac{dT(t)}{dt} = \frac{T_0 - T(t)}{\tau_T}$$

- $T_0$  is reference temperature
- $\tau_t$  is a coupling / relaxation constant
  - $\tau_t$  tiny, heat moves fast.  $\tau_t$  big, ...
- to implement this idea ? Multiply velocities

\*Slight simplification of formula

# Implementation of weak coupling

Scale velocities,  $v_{new} = \lambda v_{old}$  and  $\lambda = \left( 1 + \frac{\Delta t}{\tau_T} \left( \frac{T_0}{T} - 1 \right) \right)^{1/2}$

Intuitively

- $\Delta t$  (time step) big ? temperature will change more
- what if  $T_0 = T$  ?
- square root ?
  - wrong  $T$  reflects a difference in  $v^2$

# Importance of heat baths

Does not conserve energy

In principle

- bring a system to equilibrium for temperature

In practice

- avoid damage due to numerical errors / approximations

For a system at equilibrium

- heat bath should do nothing

Does allow artificial tricks

- gently heat a system and watch behaviour
- gently cool a system and "anneal" it (more later)

Extension to other properties

- analogous reasoning for pressure bath

# comparison of Monte Carlo and Dynamics Simulations

MC	MD
any cost/energy OK	requires continuous $E_{pot}(\mathbf{r})$
time usually no meaning	gives time scales
most moves OK	physical trajectories
temperature from acceptance/rejection	has explicit $E_{kin}$ and temperature bath
easy to program	difficult
both yield a Boltzmann distribution	
both include entropy	

# Applications MC / MD

## Basic tools

- Force field
- MD / MC

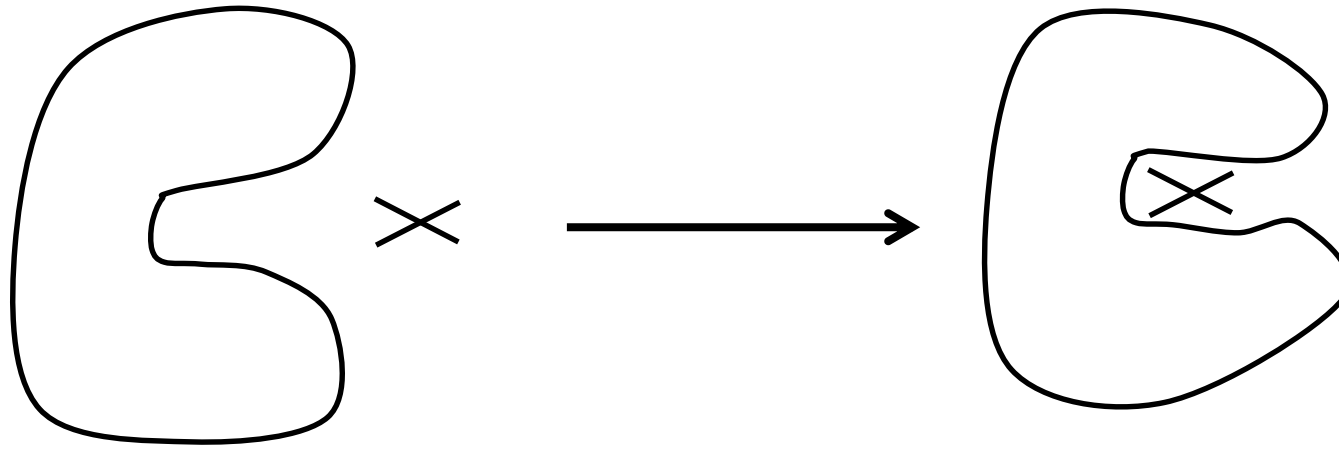
## Some application areas

- timescales
- free energy calculations
- simulated annealing
- structure refinement

# Simulating dynamics (optimistic / naïve)

## Claim

- protein has a hinge which must open to bind ligand



Can one see rates ?

- rates for different ligands ?

# Timescales

Most common quantity  $\tau$

- time to rotate by 1 radian
- time for decay in  $A(t) = A(0)e^{\frac{-t}{\tau}}$ 
  - relaxation time
  - characteristic time
- times in proteins...



# Typical times in proteins

	Amplitude (Å)	$\log_{10} \tau(\text{s})$
bond vibration	0.01 – 0.1	-14 to -13
rotation of surface sidechain	5 – 10	-11 to -10
protein hinge bending	1 – 20	-11 to -7
rotation of sidechain in middle of a protein	5	-4 to 0
local loss of protein structure	5 – 10	-5 to +1

# Timescales, simulations, statistics

Typical big simulation  $\approx 100 \text{ ns} = 10^{-7} \text{ s}$

- Imagine event with characteristic time  $10^{-7} \text{ s}$  - may or may not be seen

Consider time  $10^{-8} \text{ s}$

- may be seen a few times

What you would like - 100's or 1000's of observations

fast events

$$\tau \ll t_{simulation} \quad \text{OK}$$

$$\tau < t_{simulation} \quad \text{poor statistics}$$

slower events

$$\tau \approx t_{simulation} \quad \text{no idea / very bad statistics}$$

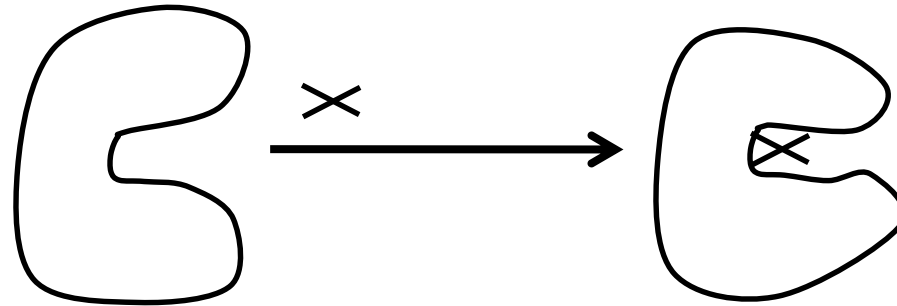
Previous example (drug binding)

- it is not enough to observe an event once (or few times)

# Free Energy Calculations

$$k_d = \frac{[\text{drug}][\text{protein}]}{[\text{drug-protein}]} = \frac{[D][P]}{[DP]}$$

$$= e^{\frac{-\Delta G}{RT}}$$



Contributing terms ?

- ligand-water  $\rightarrow$  ligand + water (many interactions,  $\Delta S$ )
- ligand+protein
- ligand loss of entropy / water entropy change
  - simulate ?

# Infinite time - free energy estimate



$$\Delta G = kT \ln \frac{[D][P]}{[DP]}$$

Very simple - simulate for long time

- Ligand (drug) goes on and off protein
- Look at [D], [P] and [DP] - calculate  $\Delta G$  directly from concentrations

Will not work - cannot simulate long enough

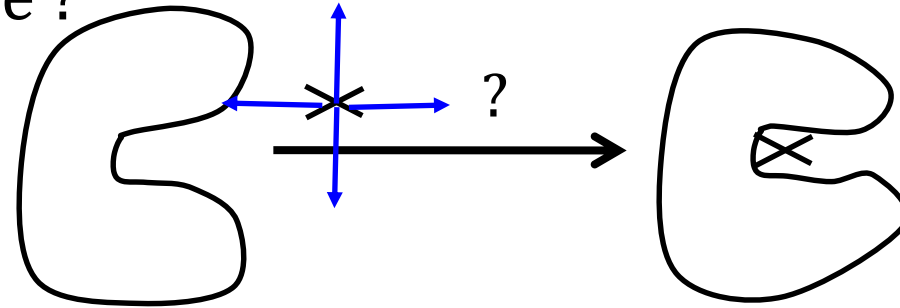
Coming philosophy

- $DP \rightleftharpoons D + P$  is too hard, find an alternative

# Free simulation for binding

If we simulate, where will the ligand go ?

What is the shape of the energy landscape ?



May take years for ligand to find protein

Short cut ?

- force ligand to protein
  - artificial force + corrections
  - very difficult – still requires rearranging water
  - entropy estimation very difficult

# Estimating free energy differences

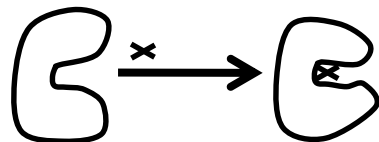
$$G = U - TS$$

$$\text{but } S = -k \sum_{i=1}^{N_{state}} p_i \ln p_i$$

- so we cannot really get  $S$
- similar problem – especially visiting high energy regions

Forget absolute free energies

- concentrate on  $\Delta G$
- no problem – usually interesting property



# Summarise free energy problem so far

- Sounds easy, just estimate  $[D]$ ,  $[P]$ ,  $[DP]$  – will not work – no simulation long enough
- Cheat – push ligand in ? System not at equilibrium, requires work
- Chemically difficult – lots of interactions
  - requires completely changing water configuration
  - breaking ligand-water interactions, finding the correct ligand-protein binding
  - big change in solvent entropy, ligand entropy, protein entropy

How can one minimise the problems ?

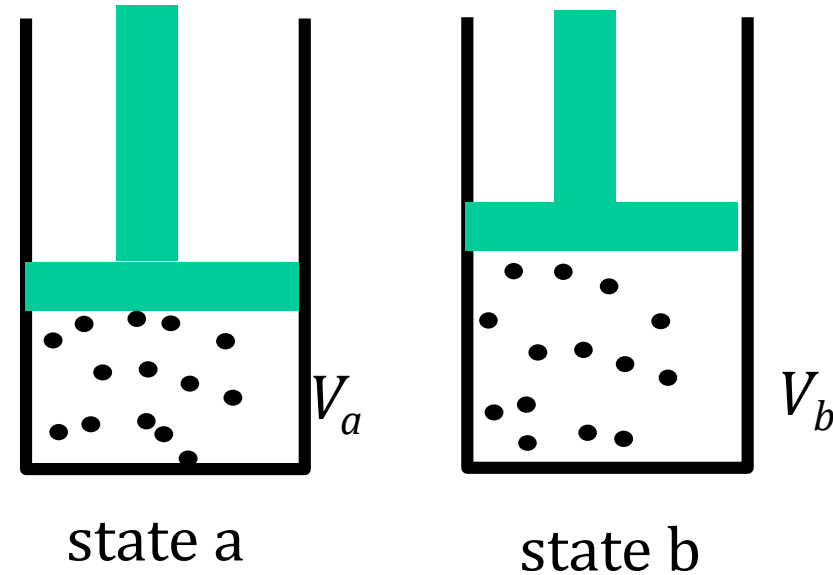
- do an easier problem (soon)

First - small detour on work

# Work and free energy changes

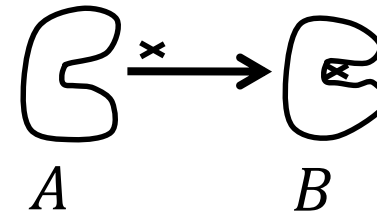
work done A to B

- free energy change
  - automatically includes entropy
  - go in either direction



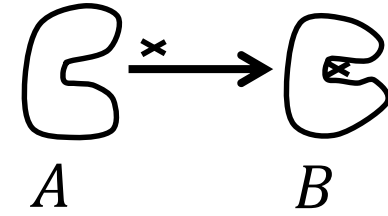
Work going from unbound  $\rightarrow$  bound

- $\Delta G_{AB}$
- what is B ? what is A ?
  - more later
- measuring work ?

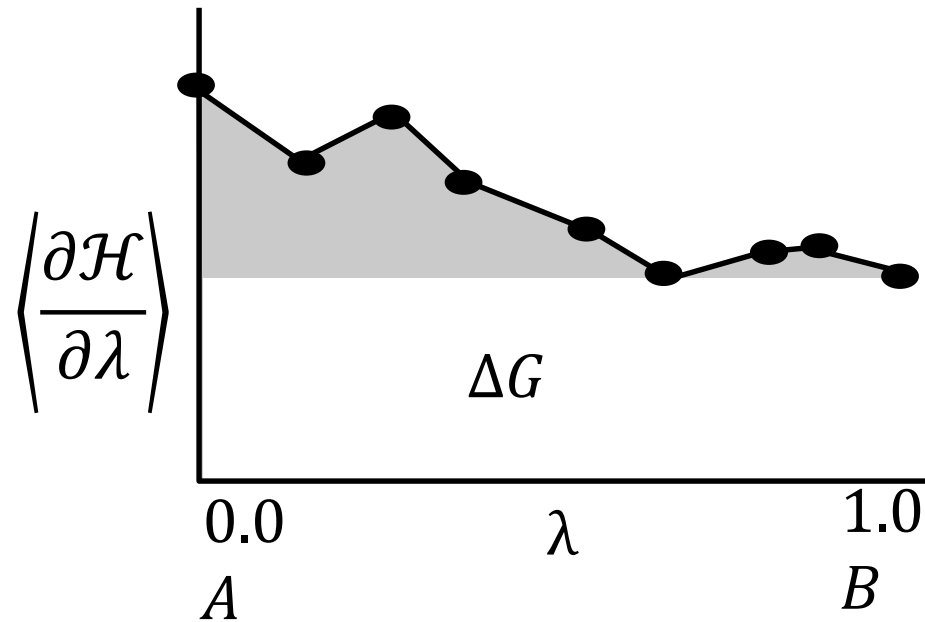




# Work and free energy



Measure the work needed to move from  $A$  to  $B$



$\mathcal{H}$  is Hamiltonian ( $E_{kin} + E_{pot}$ )

$$\Delta G = \int_A^B \left\langle \frac{\partial \mathcal{H}(\mathbf{p}, \mathbf{r})}{\partial \lambda} \right\rangle_{\lambda} d\lambda \quad \text{or} \quad \Delta G = \sum_{i=0}^{N_{step}} (H_{i+1} - H_i)$$

# Binding energy - feasibility

Would this approach work ?

$\langle \partial \mathcal{H} / \partial \lambda \rangle$  must be a good average (lots of fluctuations)  
must change  $\lambda$  slowly

Chemistry problems: your simulation would

- get averages with all water molecules
- gradually remove water molecules (high energy ?)
- find the correct binding
- get good averaging there
  
- states A and B are very different they must be well sampled
- intermediate (higher energy states) must also be sampled
- does not work well in practice

# Paths / Energy differences (detour)

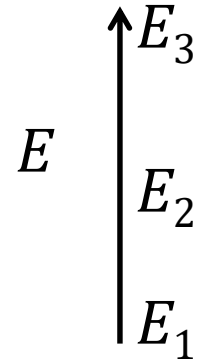
Problem – the path is too difficult – changes too big

- Energy differences depend on end states – not paths
- Look at  $\Delta E_{1,2} = E_1 - E_2$ 
  - would it matter if we go  $E_1 \rightarrow E_3 \rightarrow E_2$  ?

Can we take even stranger paths ?

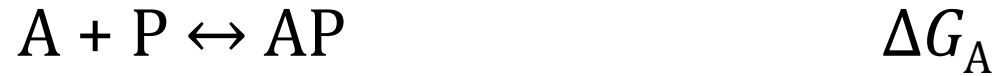
- go through non existent  $E_4$  ?
  - no problem

Same reasoning applies to free energies



# Applying different paths

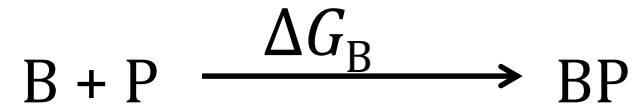
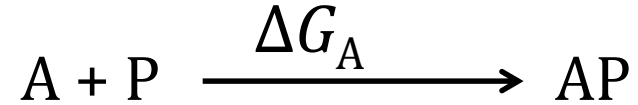
Originally wanted (ligand A or B, protein P)



If I know  $\Delta G_B$

$\Delta\Delta G_{AB}$  is easier

$$\Delta\Delta G_{AB} = \Delta G_A - \Delta G_B$$



What would  $\Delta\Delta G_{AB}$  mean ?

- relative binding strength

# Alternative routes

$\Delta G_A$  and  $\Delta G_B$  too hard

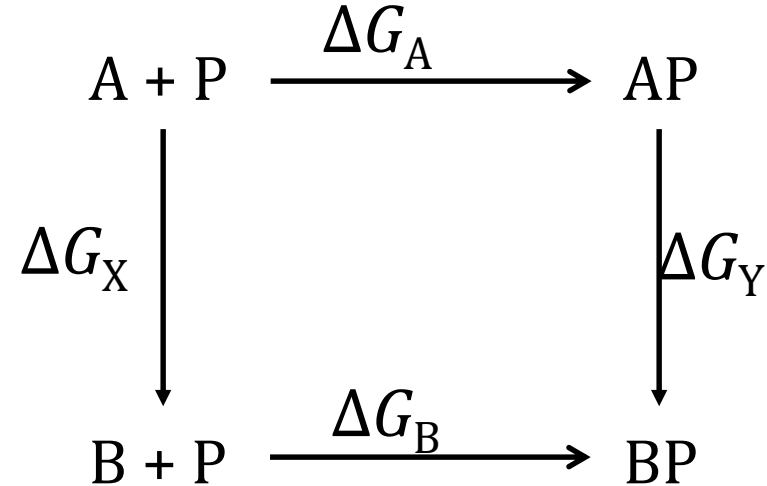
- we would be happy with  $\Delta \Delta G_{AB}$

$$\Delta G_A + \Delta G_Y = \Delta G_B + \Delta G_X$$

$$\Delta G_A - \Delta G_B = \Delta G_X - \Delta G_Y \quad \text{remember } \Delta \Delta G_{AB} = \Delta G_A - \Delta G_B$$

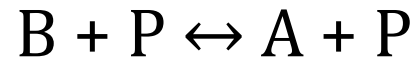
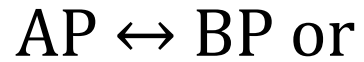
So  $\Delta \Delta G_{AB} = \Delta \Delta G_{XY}$

- why  $\Delta G_X$  easier ?
- why  $\Delta G_Y$  easier ?



# Easier free energy changes

if A/B are rather similar



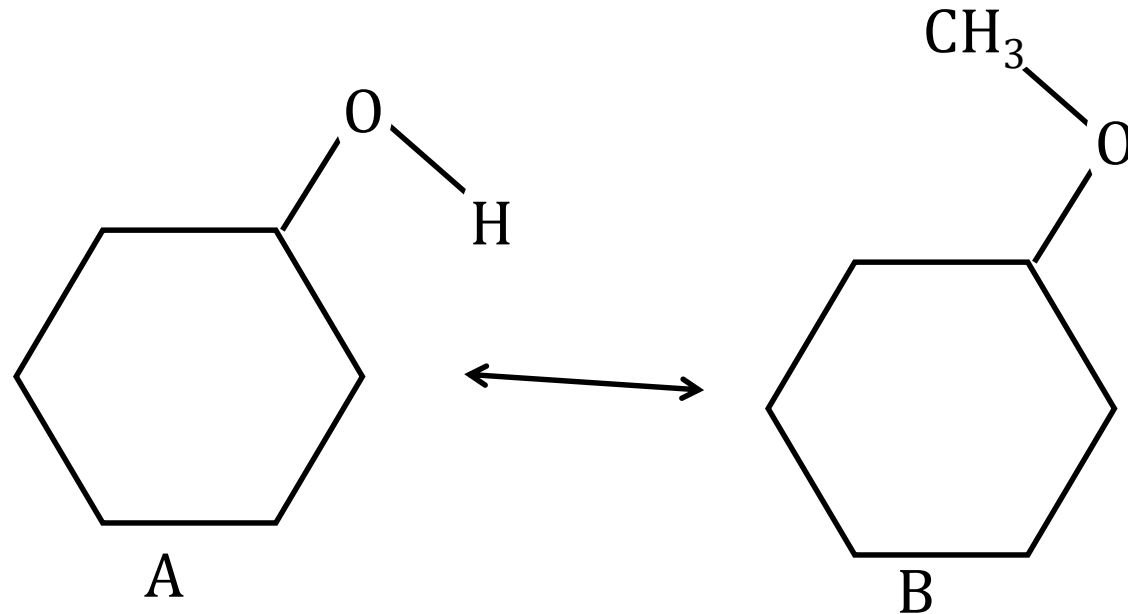
(free A  $\leftrightarrow$  B      forget the protein)

are small changes – smaller than

- removing water order, removing water energy, finding protein...

Example

- small change

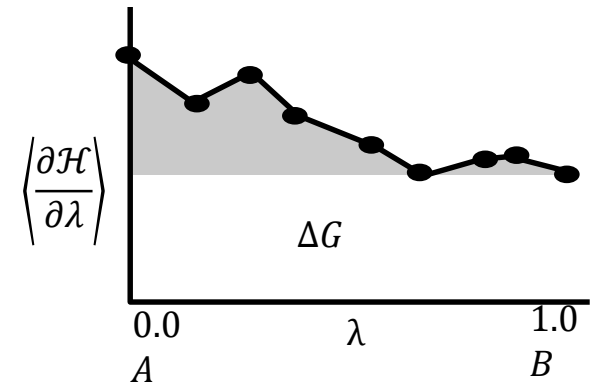
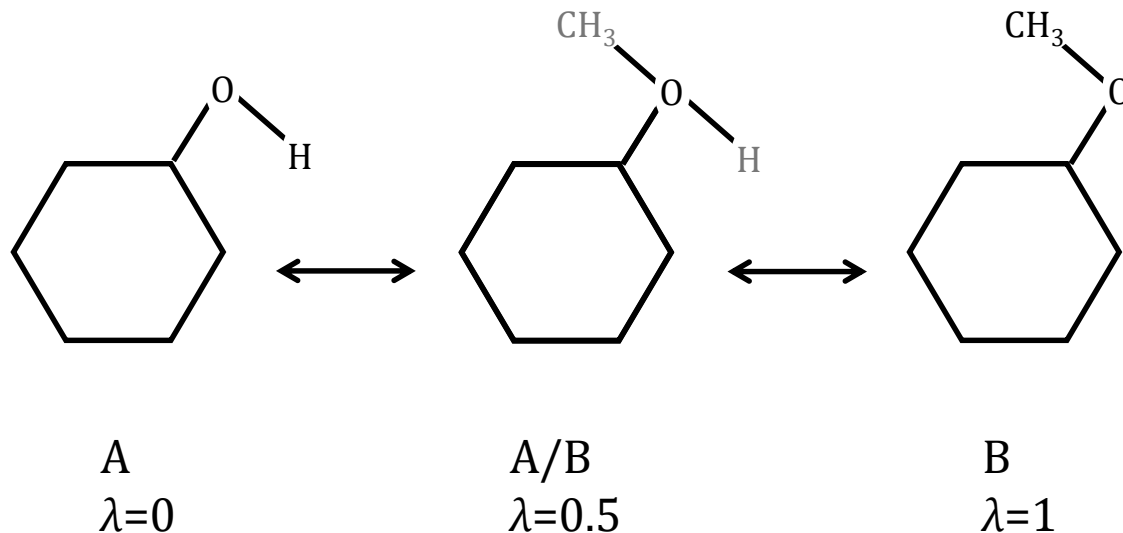


# Fictitious states

Remember formulae

$$\Delta G = \int_A^B \left\langle \frac{\partial \mathcal{H}(\mathbf{p}, \mathbf{r})}{\partial \lambda} \right\rangle_{\lambda} d\lambda \quad \text{and} \quad \Delta G = \sum_{i=0}^{N_{step}} (H_{i+1} - H_i)$$

make chemistry a function of  $\lambda$



# $\lambda$ dependence

$\lambda = 0$           an OH group

$\lambda = 1$           an OCH<sub>3</sub> group

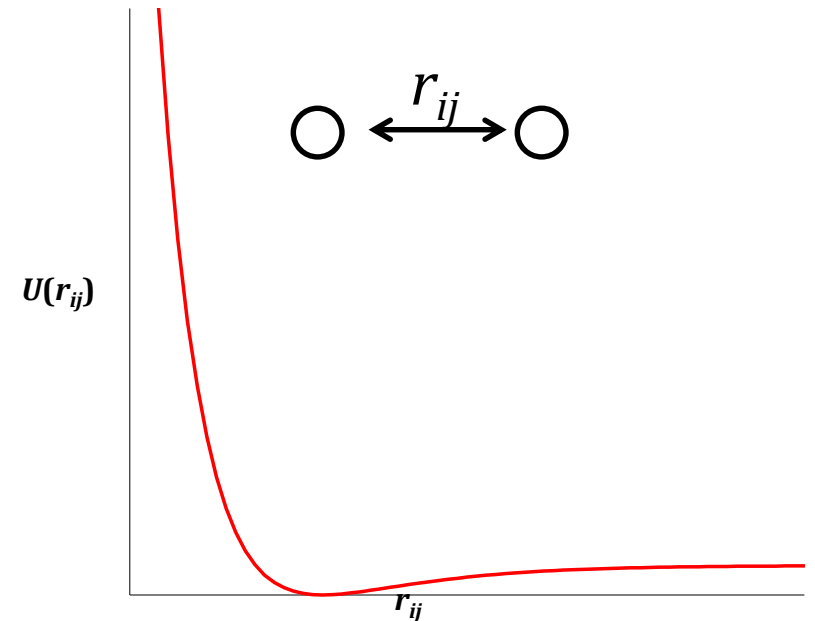
$\lambda = 0.5$

- charge of H – half of original charge
- radius / size ( $\sigma$ ,  $\epsilon$ ) half of real value and so on

Atoms gradually

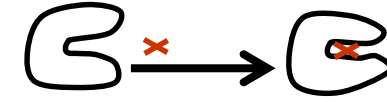
- appear in one direction
- disappear in other

Description of system is now function of  $\lambda$



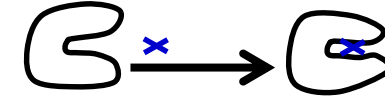


## $\lambda$ dependent simulations



Two simulations necessary

- $\lambda$  from 0.0  $\leftrightarrow$  1.0 in protein
- $\lambda$  from 0.0  $\leftrightarrow$  1.0 in water
- both from red  $\leftrightarrow$  blue



As  $\lambda$  slowly moves from 0.0

- water gradually feels more/less influence of some atoms
- system should not have to rearrange itself too much

When does method work best ?

- when changes are small
  - comparison of similar ligands in a protein

# Summary of free energy calculations

From first principles: free energy differences, equilibria

- easy to calculate
- in practice impossible (sampling not possible)

Forget absolute free energies

- $\Delta G$  determine most phenomena in the world

Processes like binding still too difficult to simulate

- slow, too many conformations / states to visit

Most calculations use  $\Delta\Delta G$

- aim to get relative binding strengths

# Simulated Annealing

Classic reference – in stine

Basic tools

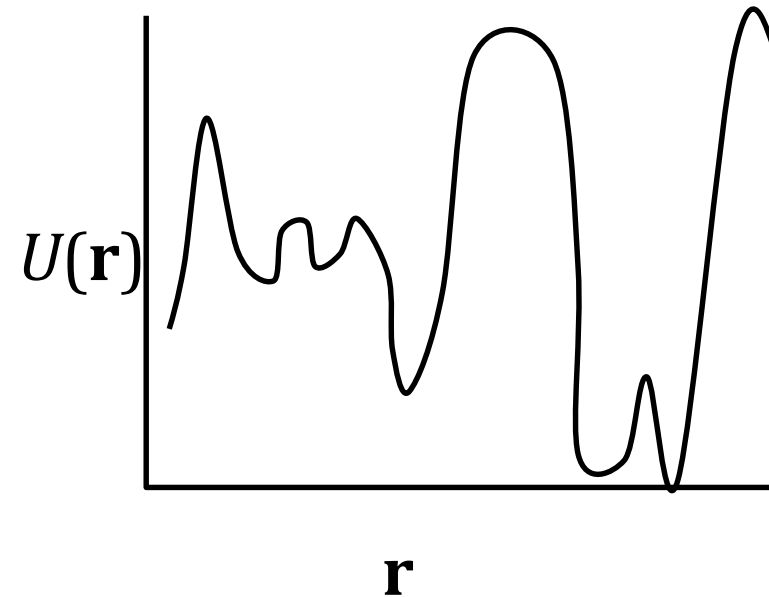
- MC or MD
  - with control of temperature (temperature bath)

Use : difficult optimisation problem

- chip layout
- travelling salesman problem
- protein structure

Optimisation problem

- several dimensional (2 to 2000)
- many local minima



# Procedure

**while** ( $T > T_{\text{end}}$ )

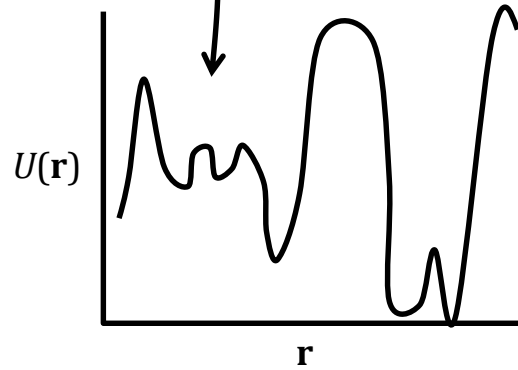
$$T(t) = T_0 e^{-ct}$$

**move system (Monte Carlo)**

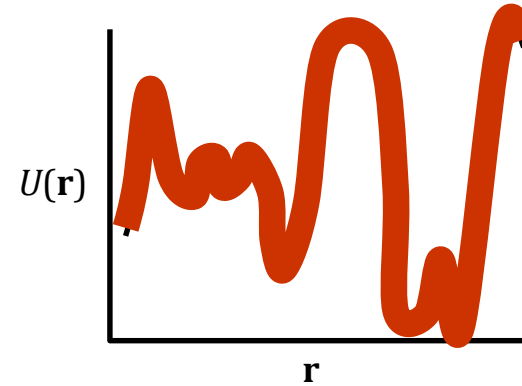
- $T_0$  initial temperature is hot
- $c$  is decay rate (cooling of system)
- cost function is
  - $E_{\text{pot}}$  in chemistry
  - path length in travelling salesman
  - board cost in chip layout problem ...
- why may this work ?

# Simulated Annealing concept

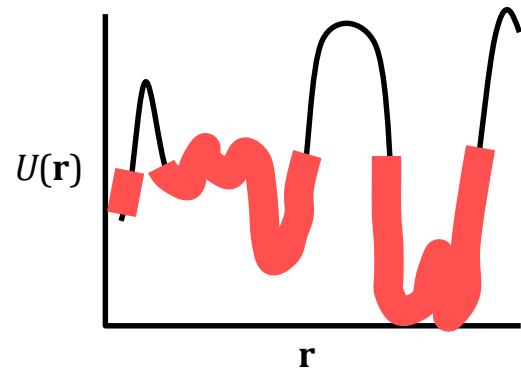
initial (poor)  
guess



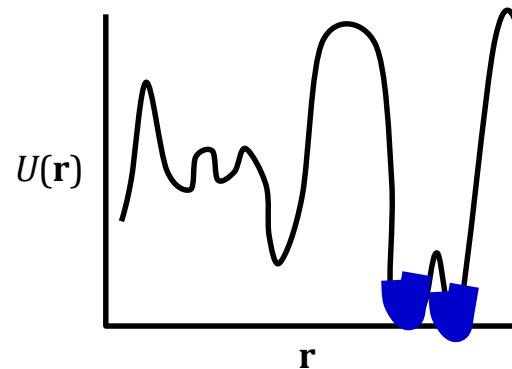
initial high  $T$   
distribution



cooler  $T$



cold  $T$



# Properties, practical issues

Admit that there may not be a best solution

- not worth spending effort between many very good solutions

Some problems have "phase transitions"

How hot should  $T_0$  be ?

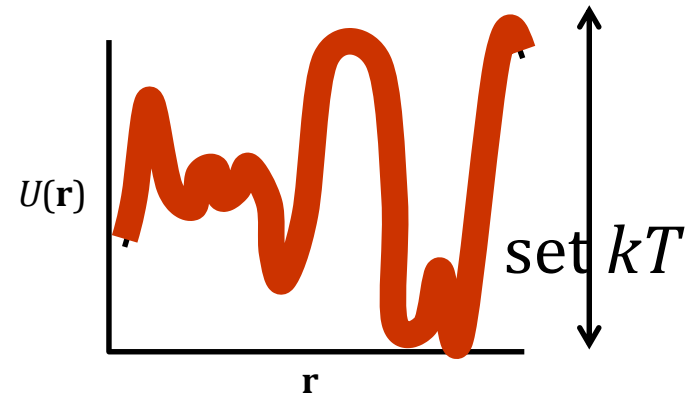
- infinite ? No : look at barriers

How slow should cooling be ( $c$ ) ?

- system should be at equilibrium
- very slow

Cool exponentially ?

- best first guess
- should certainly cool more slowly at transition points



# Anneal with MC or MD ?

Historic use of Monte Carlo

- easiest to apply to many problems

Use MD ?

- provides expected advantages (efficiency)
- uses available gradient / derivative information

Implementation

- Couple to temperature bath (dynamics) / make  $T$  time dependent (MC)

Use in practice ?

- simulated annealing in
  - most MD codes, refinement packages, ...

# Refinement of Structures (NMR / X-ray)

Story from first semester

- problem : generate protein coordinates from NMR information (or X-ray)
- distance geometry gives an initial guess, but
  - distance geometry methods spread error across all distances
  - errors are spread across bonds, measured distances
  - chirality may be broken (causes distance problems)

Belief

- coordinates are not bad, but could be improved

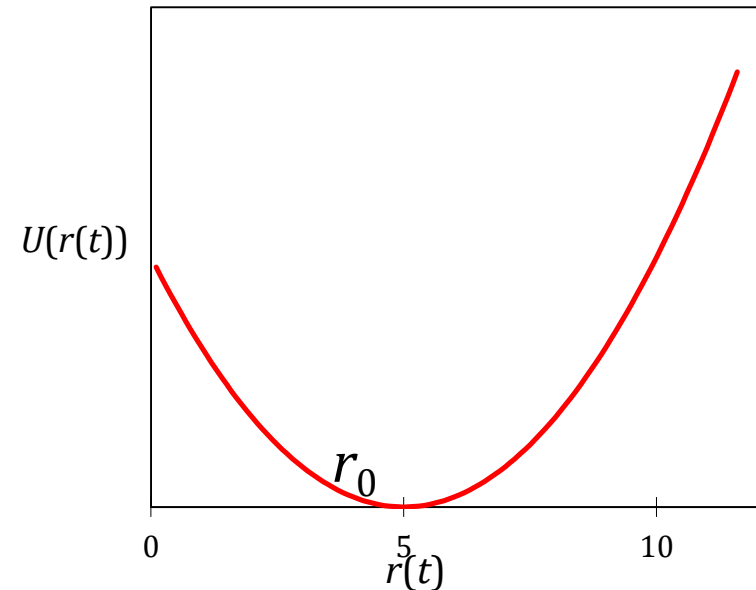


# Pseudo - energy terms

For some distance measurement  $i$  between some pair of atoms

- $r_0$  measured distance
- $r(t)$  distance between particles at time ( $t$ )
- say  $U_i(r) = c_i(r(t) - r_0)^2$
- add this to normal force field

$$U_{tot}(\mathbf{r}) = U_{phys}(\mathbf{r}) + \sum_{i=1}^{N_{restraints}} U_i(\mathbf{r})$$



$U_{phys}(\mathbf{r})$  normal force field - atomistic (bonds, electrostatics...)

## result ?

System moves to low energy + low fake energy

- gradually moves to agree with experimental data

Practical issues  $U_{tot}(\mathbf{r}) = U_{phys}(\mathbf{r}) + \sum_{i=1}^{N_{restraints}} U_i(\mathbf{r})$

$$U_i(r) = c_i (r(t) - r_0)^2$$

- big  $c$  very artificial
- small  $c$  system will be slightly biased to agree with experimental data

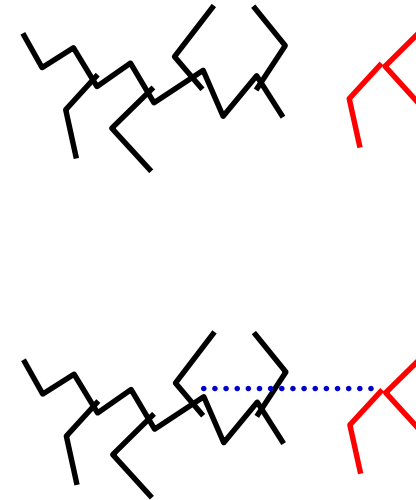
# Fake Energies - examples

Refinement of

- X-ray structures (common)
- NMR (often)
- others: microwave spectroscopy, ...

Modelling problems

- you want to put a bond in a model
  - putting it in directly
    - high energy bond
    - system stuck in minimum
  - introduce a distance restraint
    - gradually increase associated constant  $c$



# Summary

What one can do with related methods

- look at timescales of motions (very superficial)
- free energy calculations – important for problems such as binding of ligands
- simulated annealing – methods used as minimizers, not necessarily to get an ensemble
- pseudo-(potential) energies (X-ray, NMR, ...)