

Self-consistent mean field methods

General method for problems with

- multiple sites
- each site exists in different states
- each site interacts with other sites

History

- probably Ising spin model
- application different to this one

Aims

- find optimal set of states or
- find distribution of states at a given temperature

examples

Relevant to us

- protein side chains
- RNA base pairing
- sequence design

Historic / simple

- spin systems

Not here

- wave functions (standard method)
- polymer properties

Plan

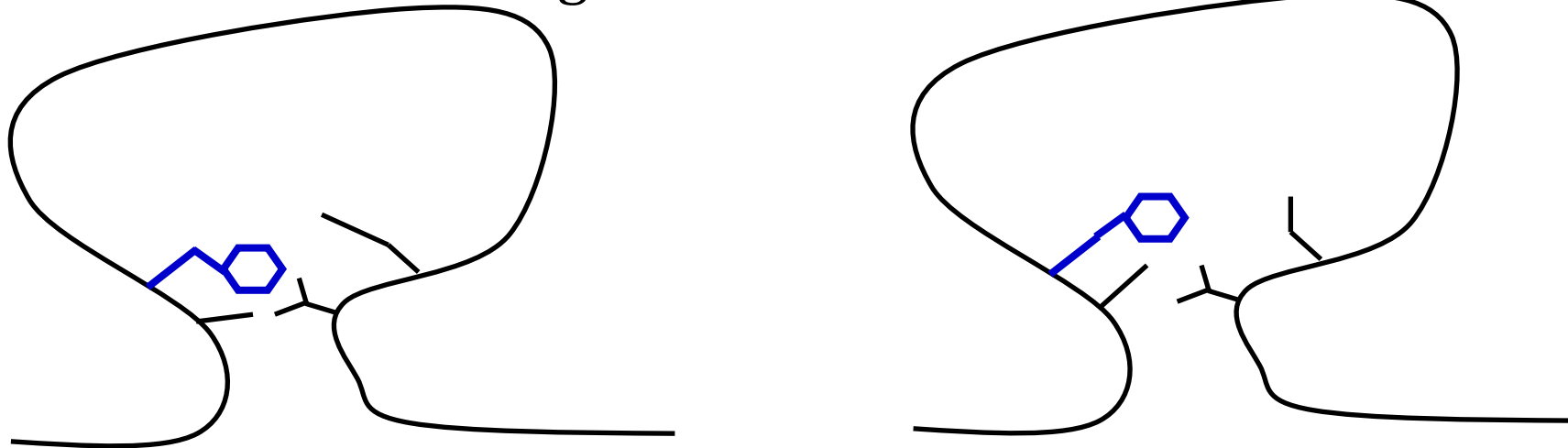
- some example problems
- Boltzmann relation
- examples in detail

Examples

- common feature
- parts of a system exist in some number of states
- parts of a system interact with each other

Protein side chains

- optimise (energy) their coordinates
- each interacts with his neighbours

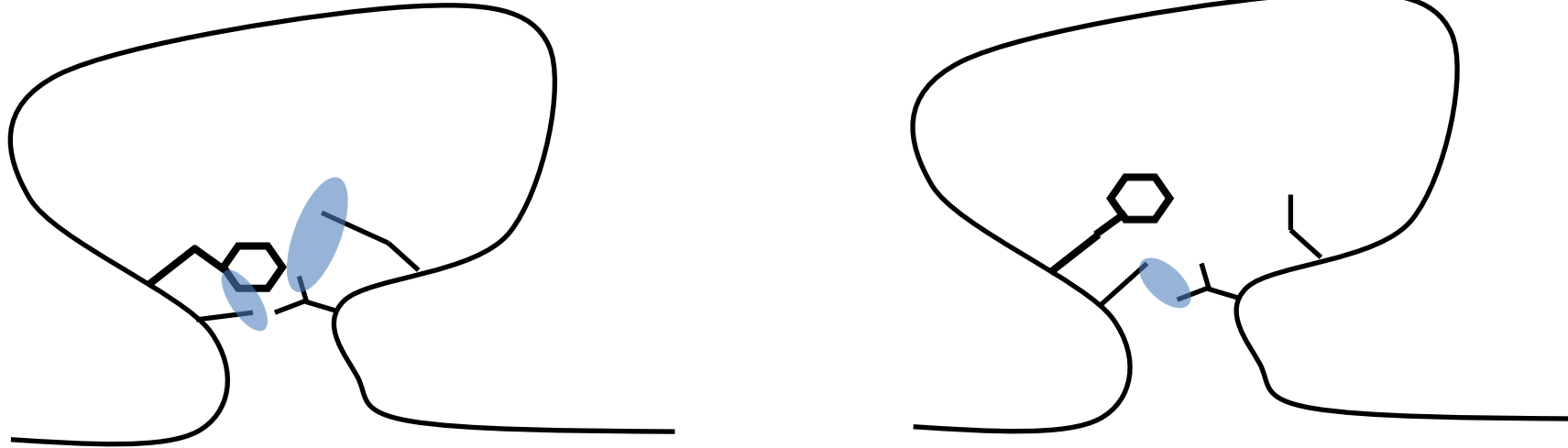


Simplification

- each sidechain can exist in one of m positions
 - say $m = 3$

Protein side chains

How many interactions ?



Make one interaction and break another

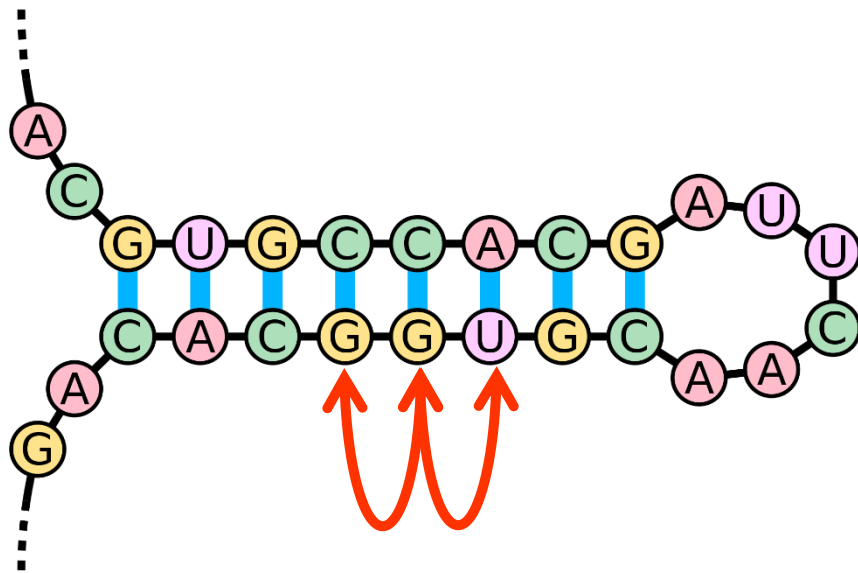
- what is the best combination ?

How big is the search space ?

- n sidechains each has m configurations = m^n
- for $m=3$ we have $3^n =$ very many

Sequence Design

RNA, but could be proteins, DNA



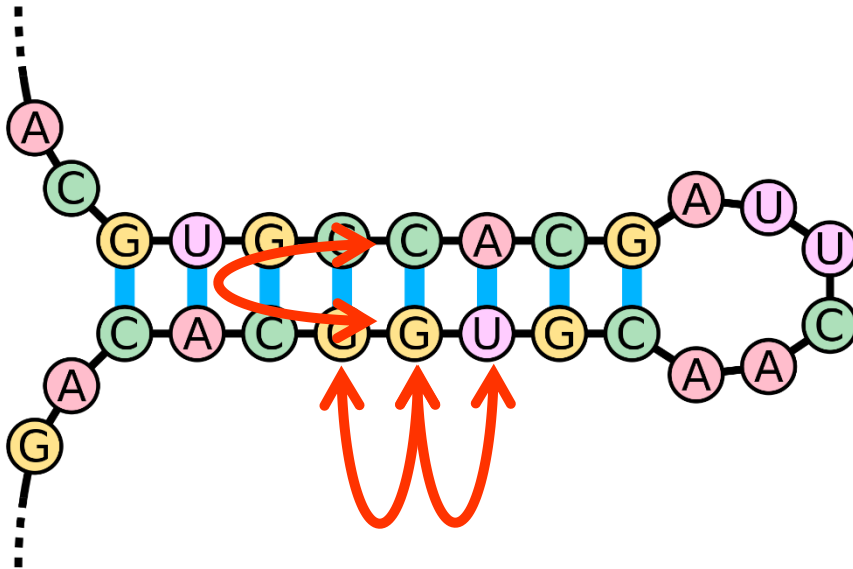
How are energies calculated ?

1. base pairs – across chain
2. sequence neighbours – base stacking

Sequence Design

Best energy

- change one base
 - affects neighbours
 - across
 - along chain



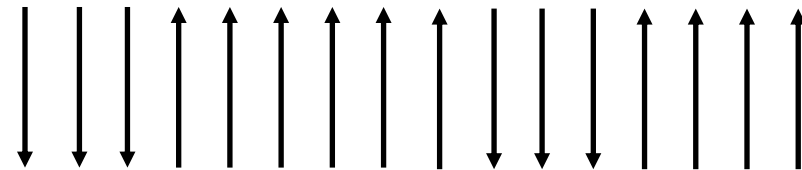
- $m = 4$ base types
- $n =$ length positions
- m^n possibilities (search space)

magnetism / spin models

Not bioinformatics ? Classic / historic

Energy (no external field)

$$V = -c \sum_{i=1}^{n-1} \sigma_i \sigma_{i+1}$$



- 2^n possible arrangements
- flip one spin to fit to left neighbour
 - might break interaction to right neighbour

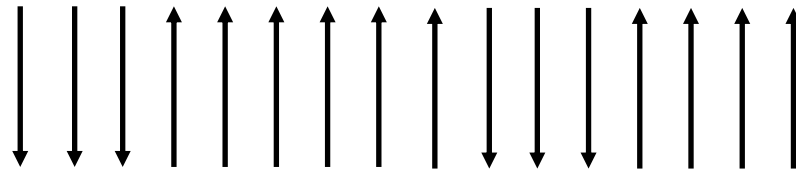
c some constant

σ_i vector – which way is spin i pointing ?

magnetism / spin models

Toy example ? You know the optimal answer(s)

Systems with more states / more complicated interactions



Do not always want the optimum

- distribution as a function of temperature

More examples

Electronic configuration of a small system (n electrons)

- shells s, p, d, \dots
- electrons have spin ($\uparrow \downarrow$)
- each electron interacts with every other electron
- put an electron in a certain p orbital
 - changes probability of neighbours
 - changing their probabilities changes

Common properties

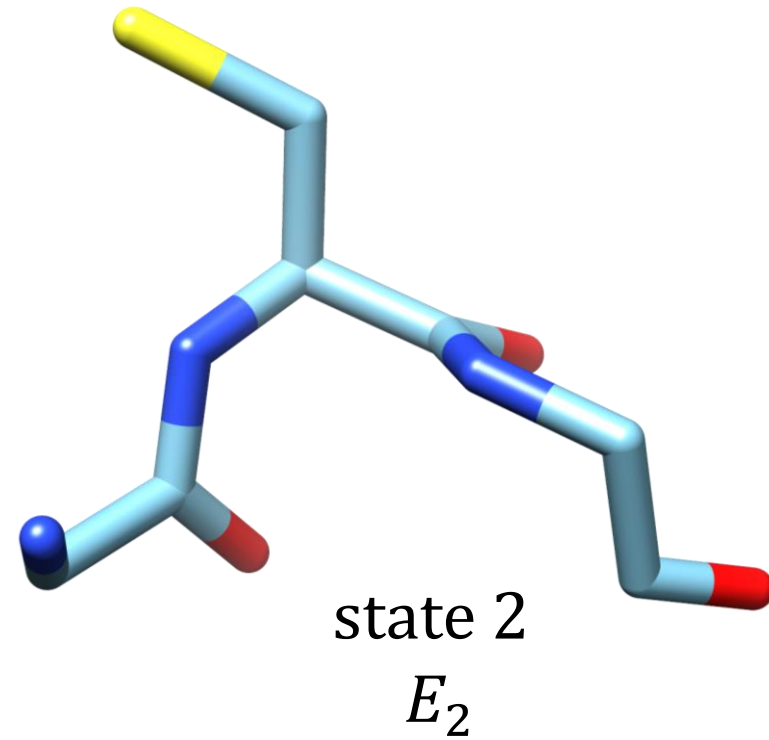
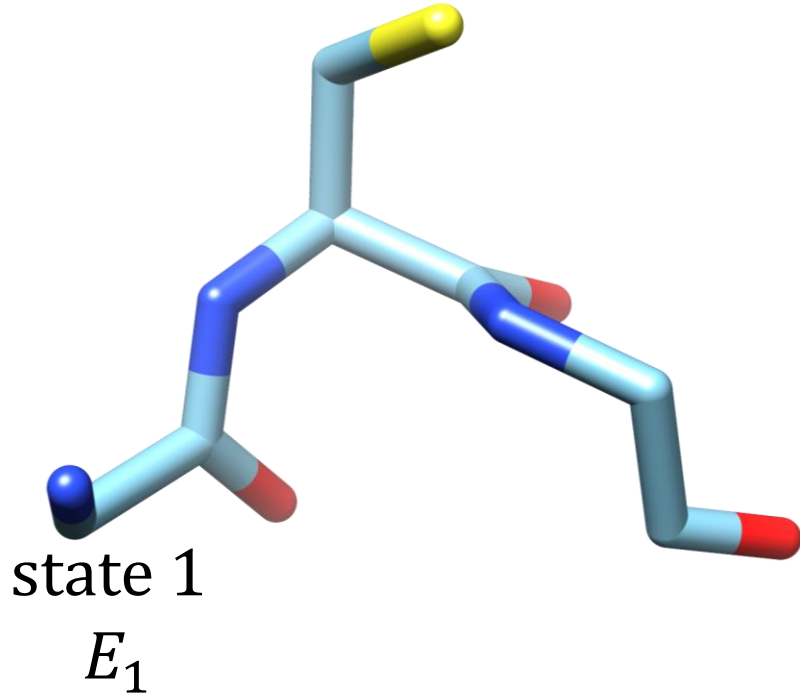
- n sites
- m states
 - m^n search space
- changing state at i affects site j which affects site k ...
- sites are not independent
 - you cannot optimise i , then j , then k , ...

General approach

- mean field methods / self-consistent mean field methods

Boltzmann .. the detour

Site with two states

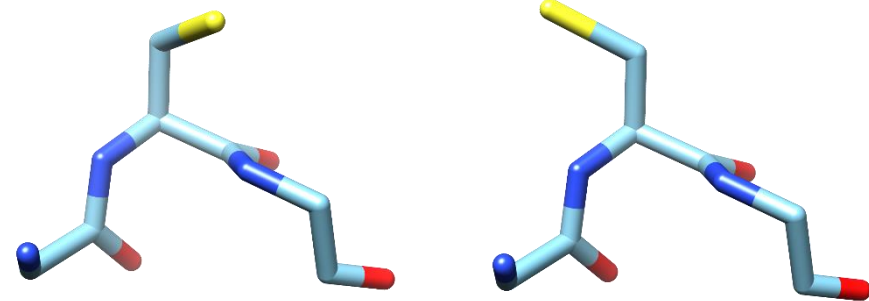


Energy difference $\Delta E = E_1 - E_2$
what is the ratio of populations ?

Boltzmann

$$\frac{p_i}{p_j} = e^{-\Delta E/kT}$$

why should you believe me ?



$$\ln \frac{p_i}{p_j} = -\frac{\Delta E}{kT}$$

$$\Delta E = -kT \ln \frac{p_i}{p_j} \text{ which looks like } \Delta G = -RT \ln \frac{[A]}{[B]}$$

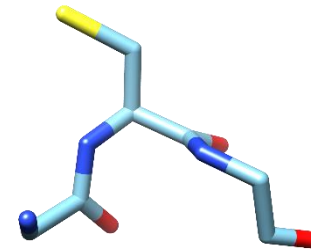
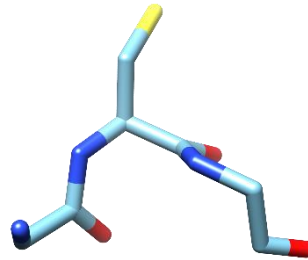
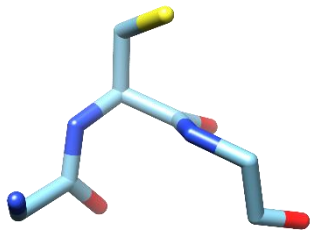
in the reaction $A \rightleftharpoons B$

Boltzmann – more states

$$\frac{p_i}{p_j} = e^{-\Delta E/kT} \quad \text{but we could also say } \frac{p_i}{p_j} = \frac{\exp\left(\frac{-E_1}{kT}\right)}{\exp\left(\frac{-E_2}{kT}\right)}$$

$e^{\frac{-E_i}{kT}}$ is the Boltzmann weight of i

What if I have three states ?



$$p_1 = \frac{\exp\left(\frac{-E_1}{kT}\right)}{\exp\left(\frac{-E_1}{kT}\right) + \exp\left(\frac{-E_2}{kT}\right) + \exp\left(\frac{-E_2}{kT}\right)}$$

what about m states ?

Boltzmann – n states

$$p_1 = \frac{\exp\left(\frac{-E_1}{kT}\right)}{\exp\left(\frac{-E_1}{kT}\right) + \exp\left(\frac{-E_2}{kT}\right) + \exp\left(\frac{-E_2}{kT}\right)}$$

generalises to

$$p_i = \frac{\exp\left(\frac{-E_i}{kT}\right)}{\sum_j^m \exp\frac{-E_j}{kT}}$$

will be used over and over again

Distributions

Simple system with two states $\frac{p_i}{p_j} = e^{-\Delta E/kT}$

At $T = 0$, $\frac{-\Delta E}{kT}$ becomes huge, negative
all the probability goes to lowest energy state

At $T \gg 0$, $\frac{-\Delta E}{kT}$ goes towards 0, $e^0 = 1$
at high temperature, $p_i \approx p_j$

For in-between ...

Optima and Distributions

$T = 0$ or $T = 300$ K or $T = 10^{10}$?

For simulations of the real world

$$T = 300 \text{ K}$$

To find the optimum

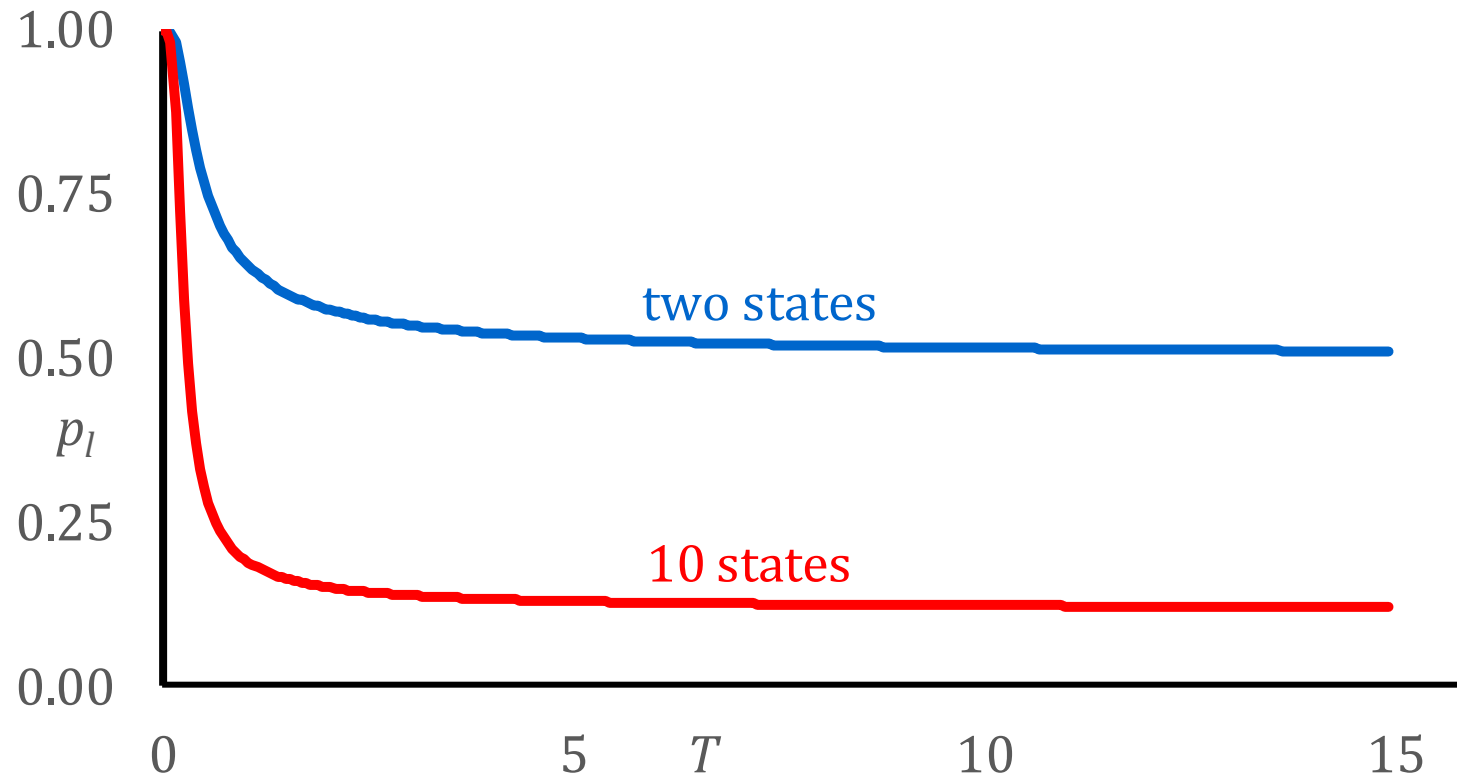
$$T = 0$$

T is

- real temperature or
- a convergence parameter
 - as the system cools, it is pushed to lower energy states

Probability as function of temperature

Probability of lowest energy state depend on T



Real world or optimisation ?

Simulations ?

- distributions of states

Answers

- rotamer distributions
- base-pairing
- sequence design
 - just the optimum

For these lectures

- best solution at $T = 0$

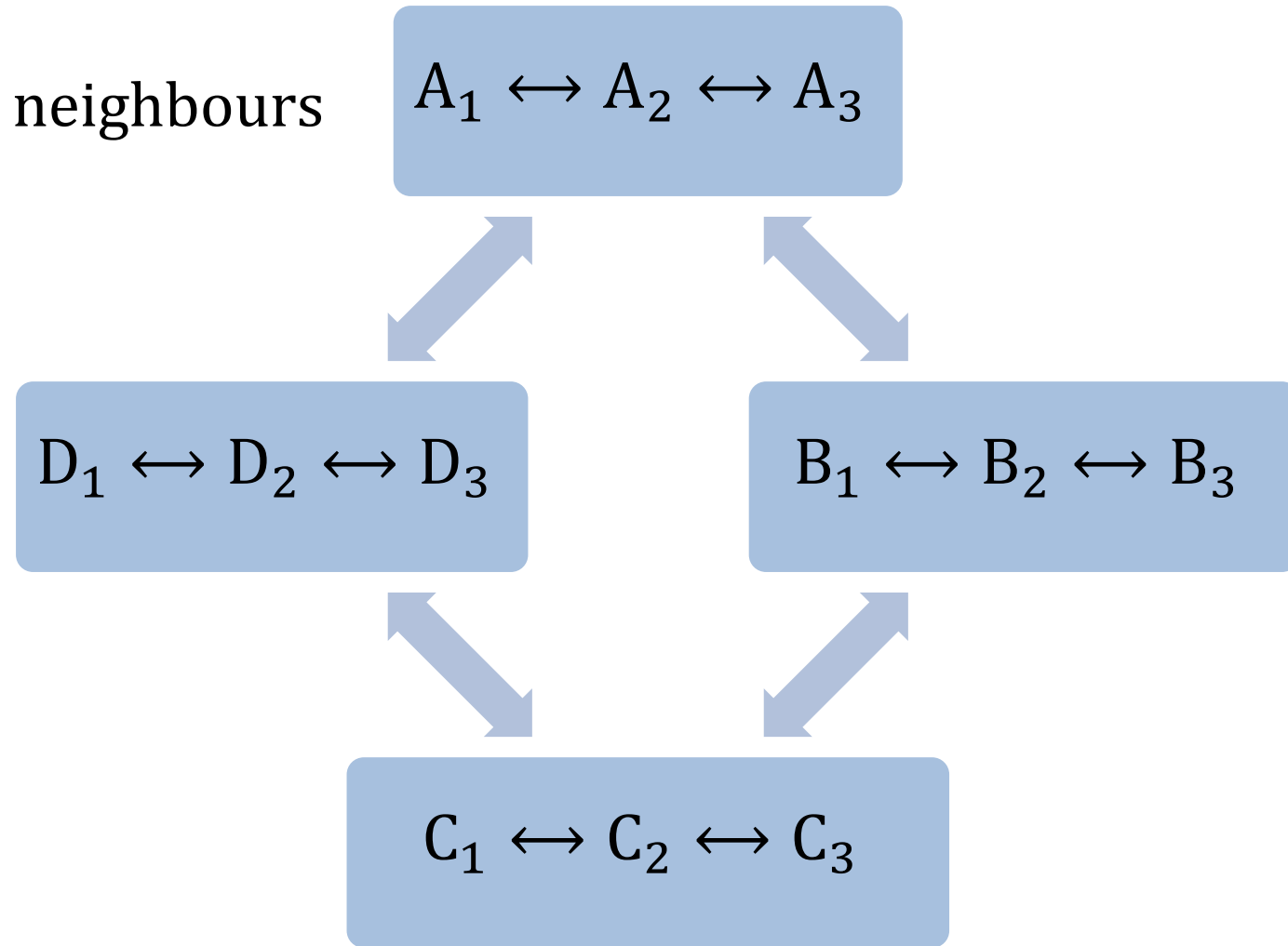
Philosophy

- start system at high temperature
 - all states are equally likely
 - each part of the system feels the average of its neighbours
- gradually cool
 - each site moves to lowest energy states

Can we just look at lowest energy state in one step ?

- no

Each site affects his neighbours

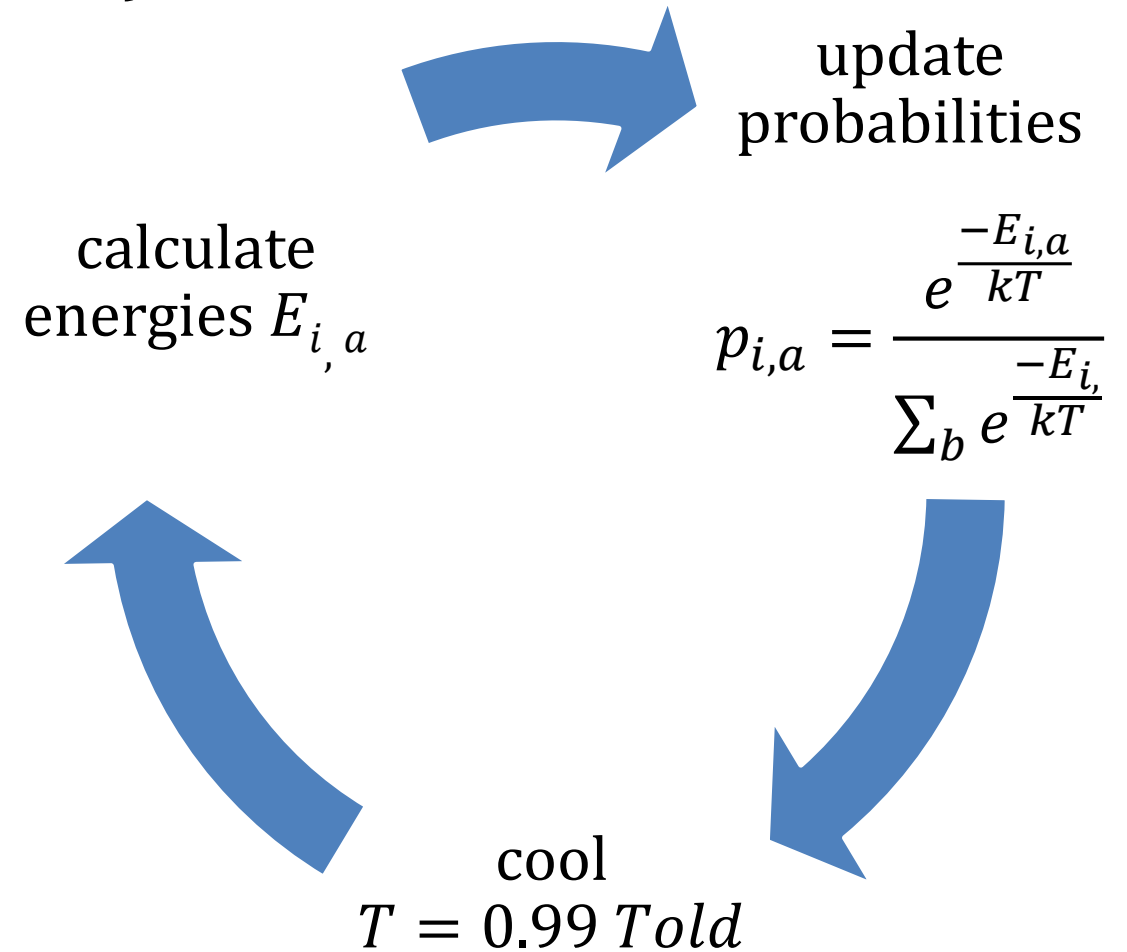


Cannot know the optimum of A since we do not yet know B, C or D

The state of A affects B affects ...

We have much bigger networks (many sites)

- adjust one a little bit, cool a bit...



Examples – Sidechain conformations

Assume

- some model for energy
- discretisation – sidechain rotamers
 - residue i can exist in m conformations

Energy depends on

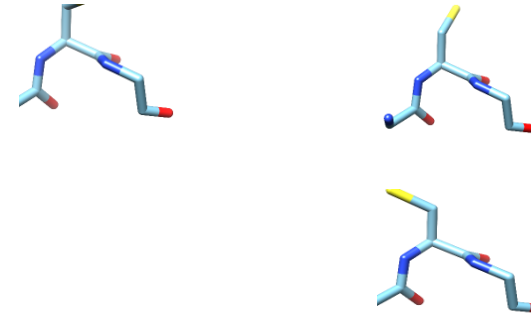
- neighbours
- interactions with backbone

- Example .. side chain with 3 positions

Work through a calculation

Use a, b, c for states... Use $i, j, ..$ for sites

Consider one side chain at site i



- 3 states (for example)
- we want probability $p_{i,a}$ in each state a

What are the interactions of sidechain i ? Consider neighbour j

- j has a probability $p_{j,a}$ of being in state a (for all the different a)
- use the mean field

mean field

Say $E(i, j)$ is the energy of sites i and j interacting, but be more specific

$E(i_a, j_b)$ is the energy due to i in state a with j in state b

We do not know the state of j , but we do know the probabilities

$$E(i_a, j) = \sum_b^{m_{states}} (p_{j,b} E(i_a, j_b))$$

this is for one neighbour, but we want the total energy $E_{i,a}$

$$E_{i,a} = \sum_j^{n_{neighbour}} \left(\sum_b^{m_{states}} (p_{j,b} E(i_a, j_b)) \right)$$

Summation over all states of neighbours – mean field

Now have $E_{i,a}$

- repeat for each state a use the Boltzmann rule to get the probabilities

- from $p_{i,a} = \frac{\exp\left(\frac{-E_a}{kT}\right)}{\sum_{b=1}^{N_{states}} \exp\left(\frac{-E_b}{kT}\right)}$

In words...

for each site i

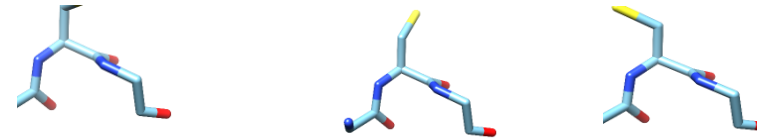
for each state a

for each neighbour j

for each state b

add in $p_{j,b} E(i_a, j_b)$

$$E_{i,a} = E_{i,a} + p(j, b)E(i_a, j_b)$$



still not finished

Why cool ?

Remember $\Delta G = -RT \ln \frac{[A]}{[B]}$ in the reaction $A \rightleftharpoons B$ so $\frac{[A]}{[B]} = e^{-\frac{\Delta G}{RT}}$

- if $T \neq 0$, we get an equilibrium, not an answer
- reason for..

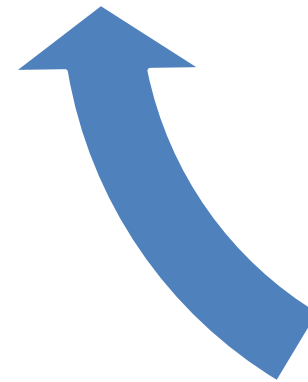


update probabilities

calculate energies $E_{i,a}$

$$p_{i,a} = \frac{e^{-\frac{E_{i,a}}{kT}}}{\sum_b e^{-\frac{E_{i,b}}{kT}}}$$

Not quite finished - initialisation



cool
 $T = 0.99 T_{old}$



Starting a calculation

- calculating $E_{i,a}$ requires knowing $p_{j,b}$ for each site j in each state b
- at the start, set all $p_{j,b}$ to $1/m_{state}$

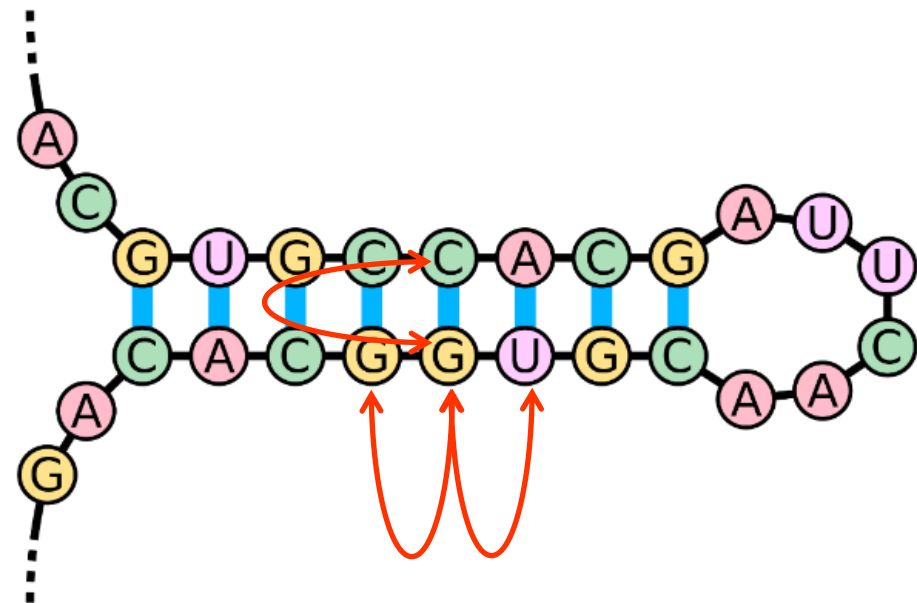
Sequence optimisation - Another example

The question

- Given a structure, find a better sequence for it
 - some energy / scoring scheme

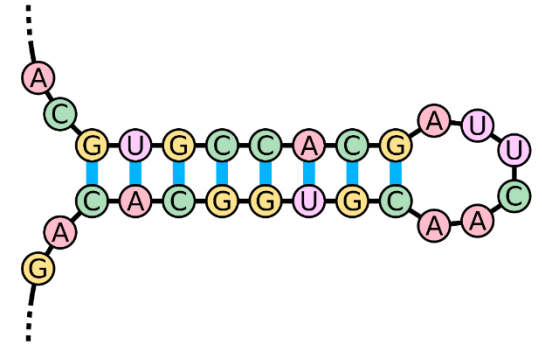
RNA or protein ? RNA in 2D is easier – energies dominated by

- base pairing
- neighbours of i ($i - 1$ and $i + 1$)



Sequence design – the philosophy

- sequence length n
- each site i
 - can be in one of four states (A, C, G, U)
 - each of the four states has equal probability $p_{i,a} = 1/n_{state} = 1/4$
- for each site i in sequence, calculate energy interacting with neighbours in each state
 - for one neighbour



$$E(i_a, j) = \sum_b^{m_{states}} (p_{j,b} E(i_a, j_b))$$

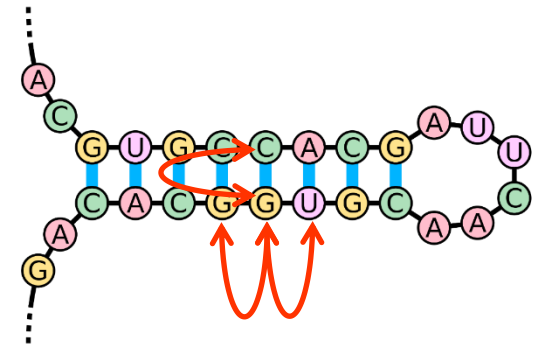
- where the summation runs over m_{states} (A, C, G, U)

- for all neighbours

$$E_{i,a} = \sum_j^{n_{neighbour}} \left(\sum_b^{m_{states}} (p_{j,b} E(i_a, j_b)) \right)$$

- neighbours are very clear here..
- then probabilities of states, at each site i

$$p_{i,a} = \frac{\exp\left(\frac{-E_{i,a}}{kT}\right)}{\sum_b \exp\left(\frac{-E_{i,b}}{kT}\right)}$$



- look at loops explicitly..

loops for sequence optimisation

for each i

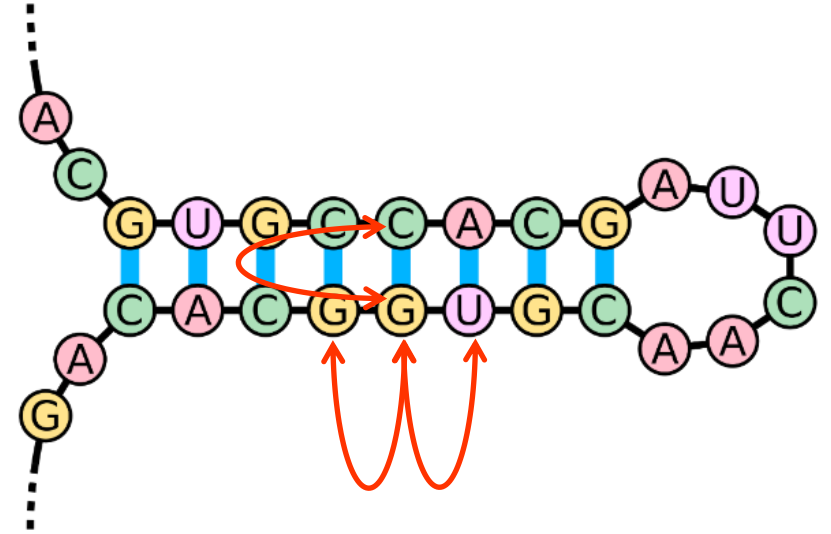
for each state a of i

for each neighbour j

for each state b of j

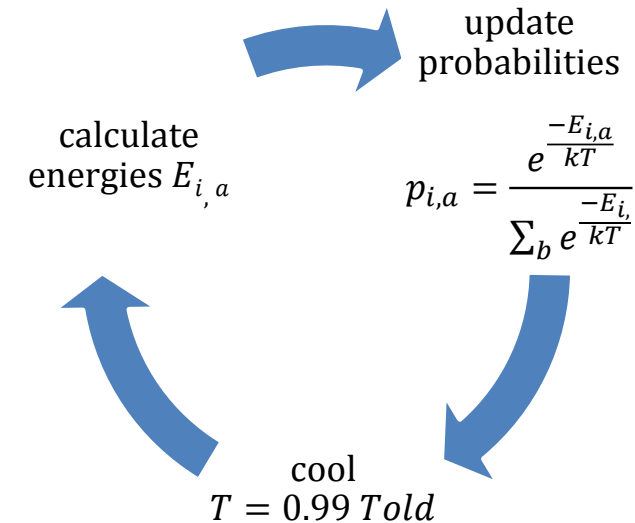
calculate $p_{j,b} E(i_a, j_b)$

store $E_{i,a} = E_{i,a} + p_{j,b} E(i_a, j_b)$



Conceptually

- at the start probabilities are all equal – no sequence
- sequence emerges at $T \rightarrow 0$



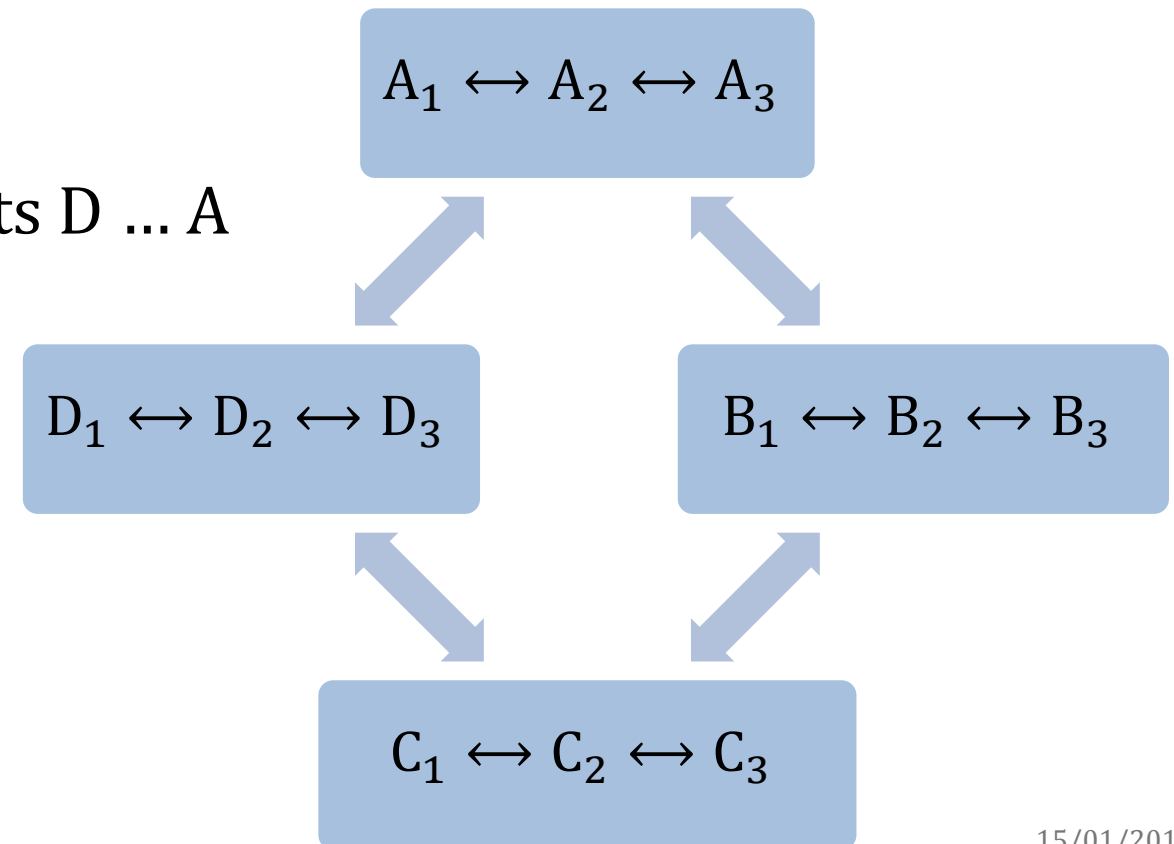
Practicalities – problems - cooling

How fast does one cool ?

- $T_{t+\delta t} = 0.99 T_t$? No. Just an example
- as in simulated annealing, cool as slowly as necessary

How slow ? Remember

- State at A affects B affects C affects D ... A
- cool slowly enough to let changes propagate / diffuse



Practicalities

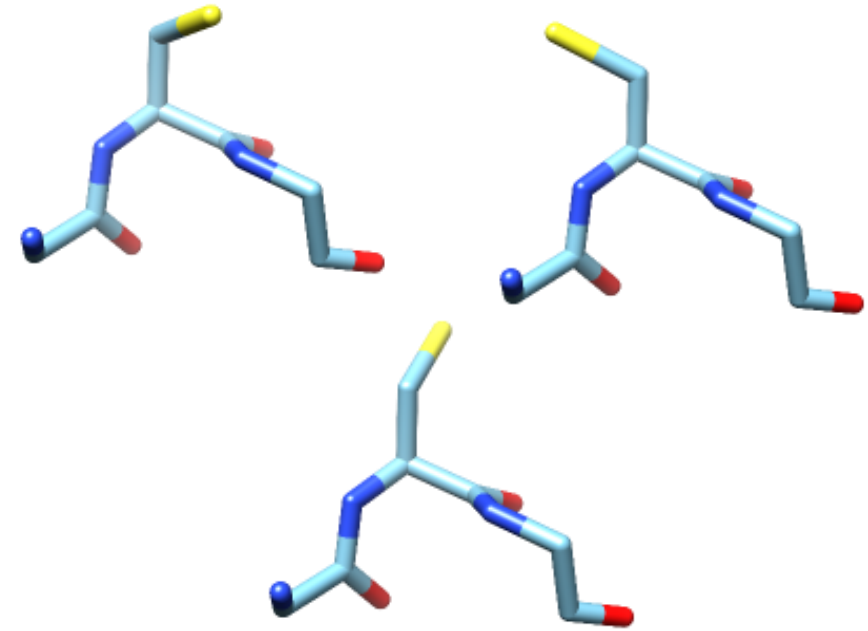
Is T a real temperature ?

- here .. No.
- in some problems.. could be

Convergence – guaranteed ?

- no

Symmetries



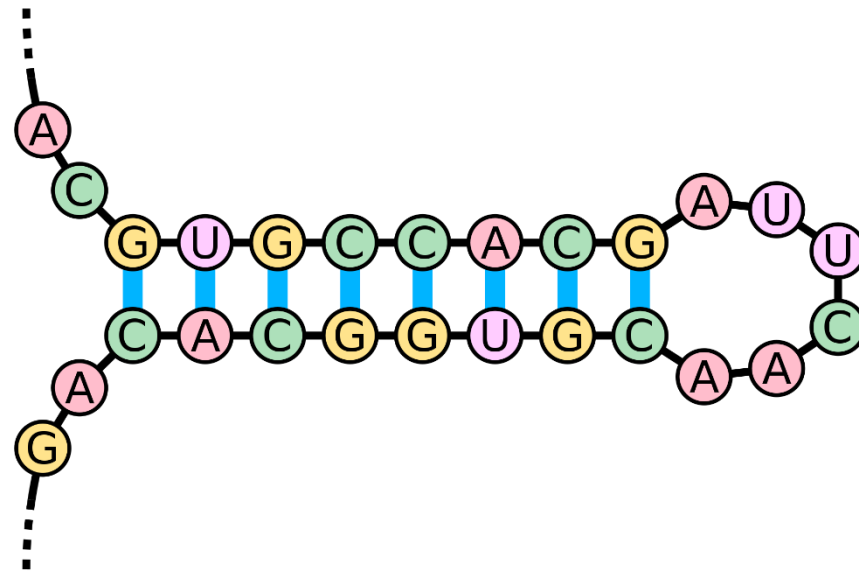
What if all rotamers have the same energy ?

- will it happen ? No
- think of interactions with backbone
 - some E_a will be better than others

RNA problem

- not as symmetric as it looks ?

What if it happens anyway ?
(same energies for different states)



Symmetries

Symmetry problems

$$p_{i,a} = \frac{1}{n_a}$$

but use

$$p_{i,a} = \frac{1}{n_a} \pm \delta \quad \text{for some very small } \delta$$

- this is enough to make one solution preferred and dominate

Lots more problems – not here

- cooling, phase transitions, oscillations, ..

Convergence

While cooling, monitor convergence

- how frozen is a system ?

In these systems, easy to measure

- at any site, can measure entropy

$$S = - \sum_{a=1}^{m_{states}} p_a \log p_a$$

maximum entropy if all states equal

$$S_{max} = -m p_a \log p_a = -m \frac{1}{m} \log \frac{1}{m} = \log m$$

minimum

$$S_{min} = 0$$

Convergence

entropy

- sum over all positions
- for fun – consider $\log_m p_a$ - gives nice normalisation

what else can one do ?

- RNA base-pairing
- sequence alignments (very difficult)
- graph problems – knapsack, bipartition

- time to stop

Summarise

Easiest on systems

- which can be discretised
- probabilities can be calculated

Works well on

- systems with many interacting parts – too big to tackle by other means
- lots of graph problems

Philosophy

- system visits all states at start
- cooling persuades it to find an answer